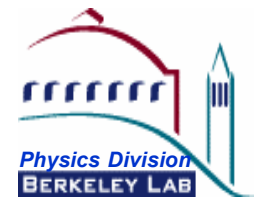


Gauge-boson Physics with ATLAS at the LHC

*with an emphasis on:
Triple Gauge-boson Couplings and
Monte Carlo Techniques for QCD Corrections*

Matt Dobbs

*Lawrence Berkeley Laboratory, USA
(U. Victoria, Canada)*



Outline



- LHC & the ATLAS Detector

- Performance example: Hadronic Endcap Calorimeter

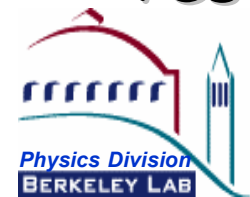
- Gauge-boson Physics

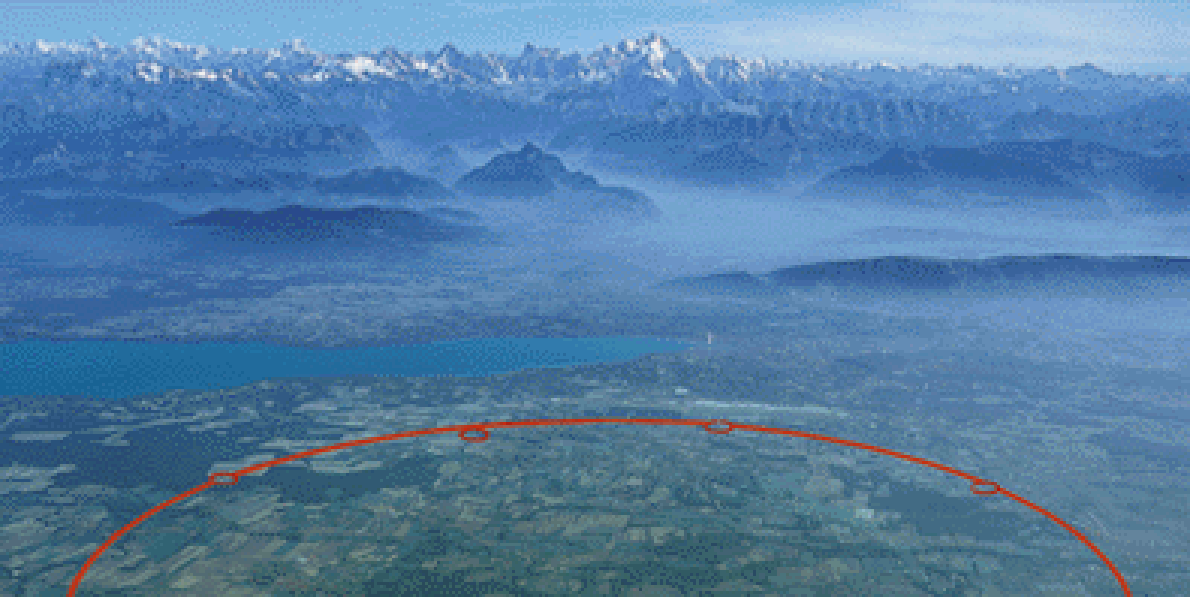
- Measuring A_{FB} and $\sin^2\theta_W$

- (Modeling our Predictions:

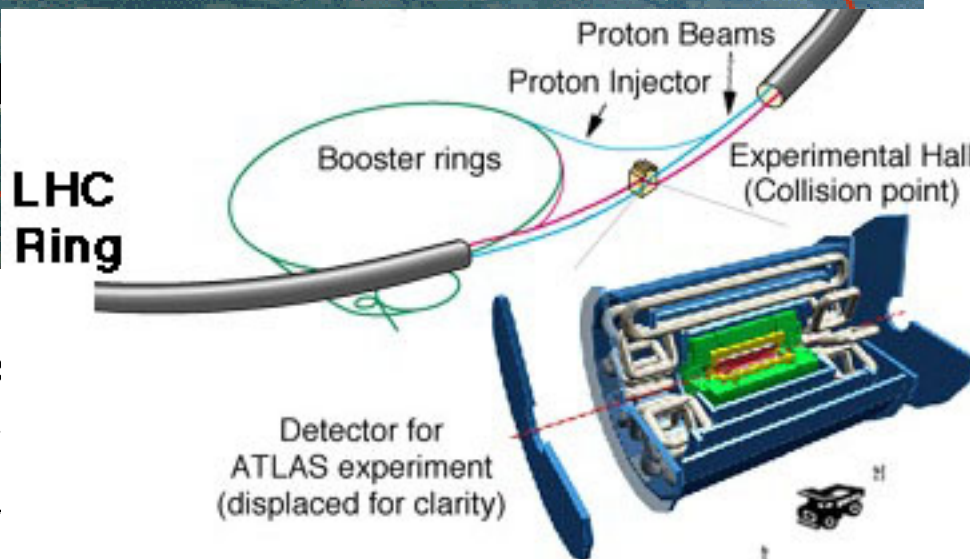
- New Monte Carlo Techniques for combining NLO(α_s) matrix elements with the parton shower
- sketch of problem to be solved, results & implications (but no details)

-) Testing the SM with Triple Gauge-boson Couplings





Large



q & g collider,

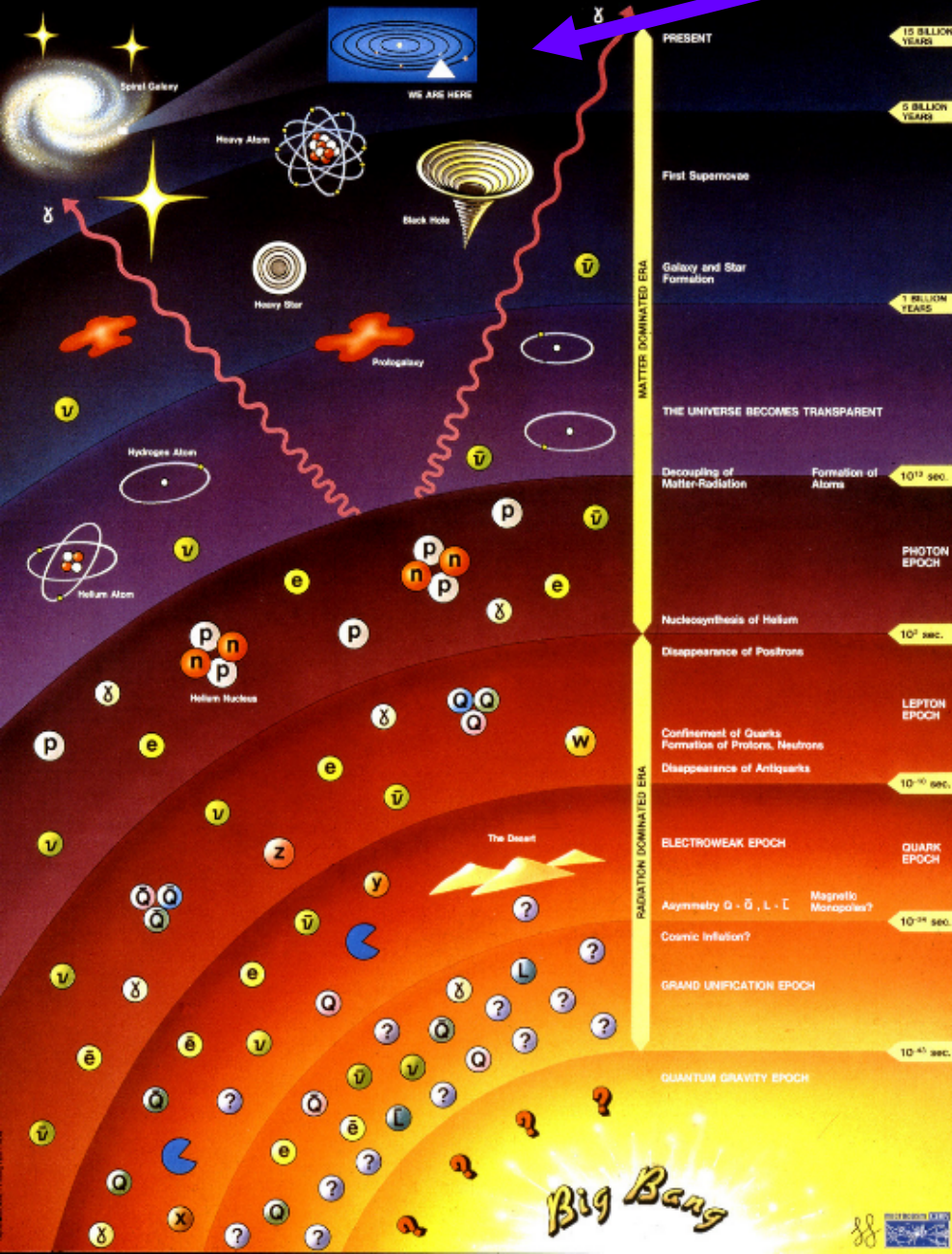
fb^{-1} in ≤ 10 years

(25ns bunch spacing)

- 14 TeV
 - ⊗ scale
- Low L \rightarrow
- High L -
 - ⊗ 25 interactions per bunch crossing
 - ⊗ 40 MHz crossing rate at high L

History of the universe

You are here



← NOW (15 Billion years)

← Stars form (1 Billion years)

← Cosmic Microwave Background
← Atoms Form (300 000 years)

← Nuclei Form (180 seconds)

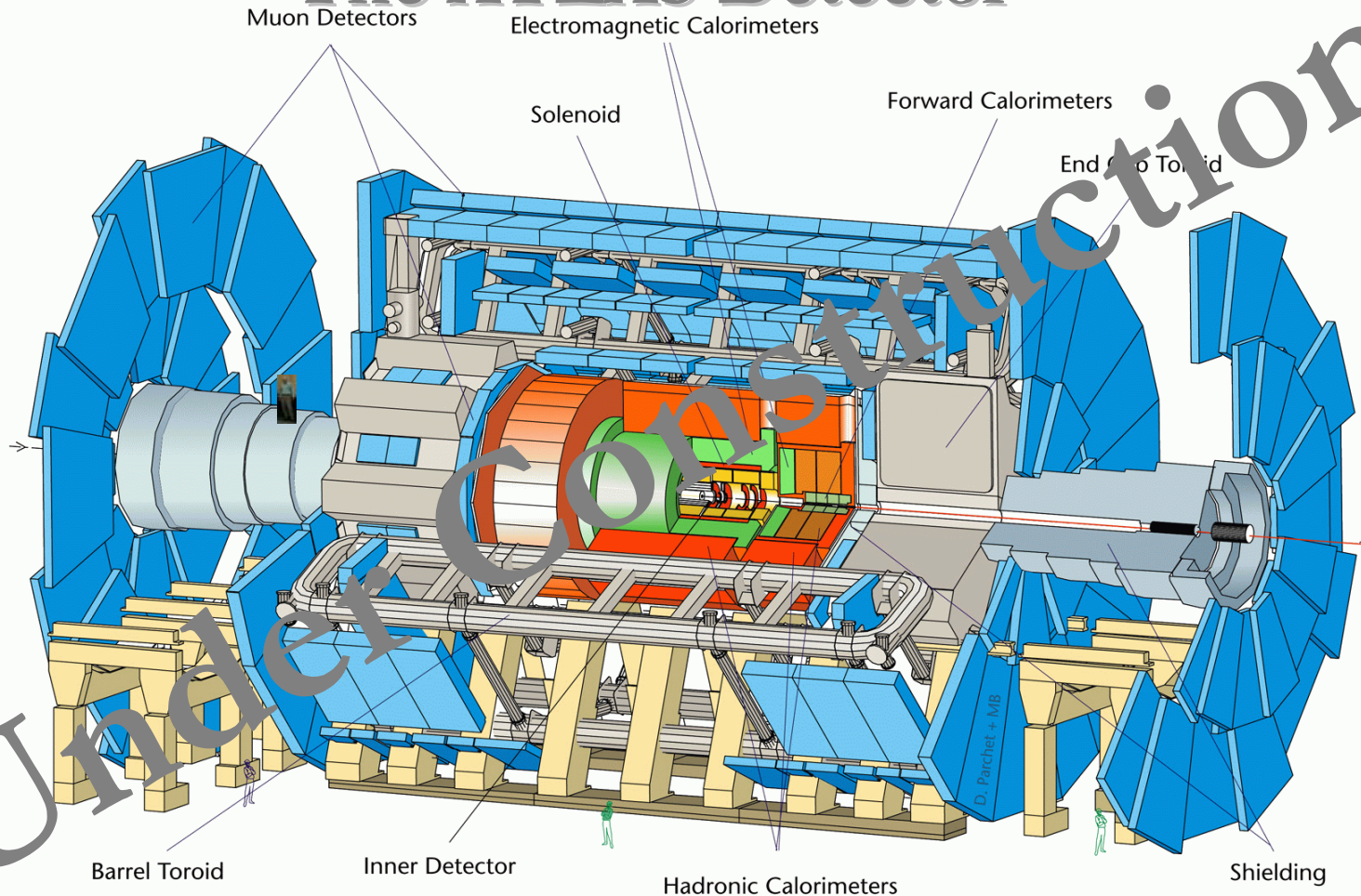
← Protons and Neutrons Form (10^{-10} sec)

← Quarks Differentiate (10^{-34} sec ?)

LHC probes physics
relevant to the universe
at age 10^{-14} sec.



The ATLAS Detector



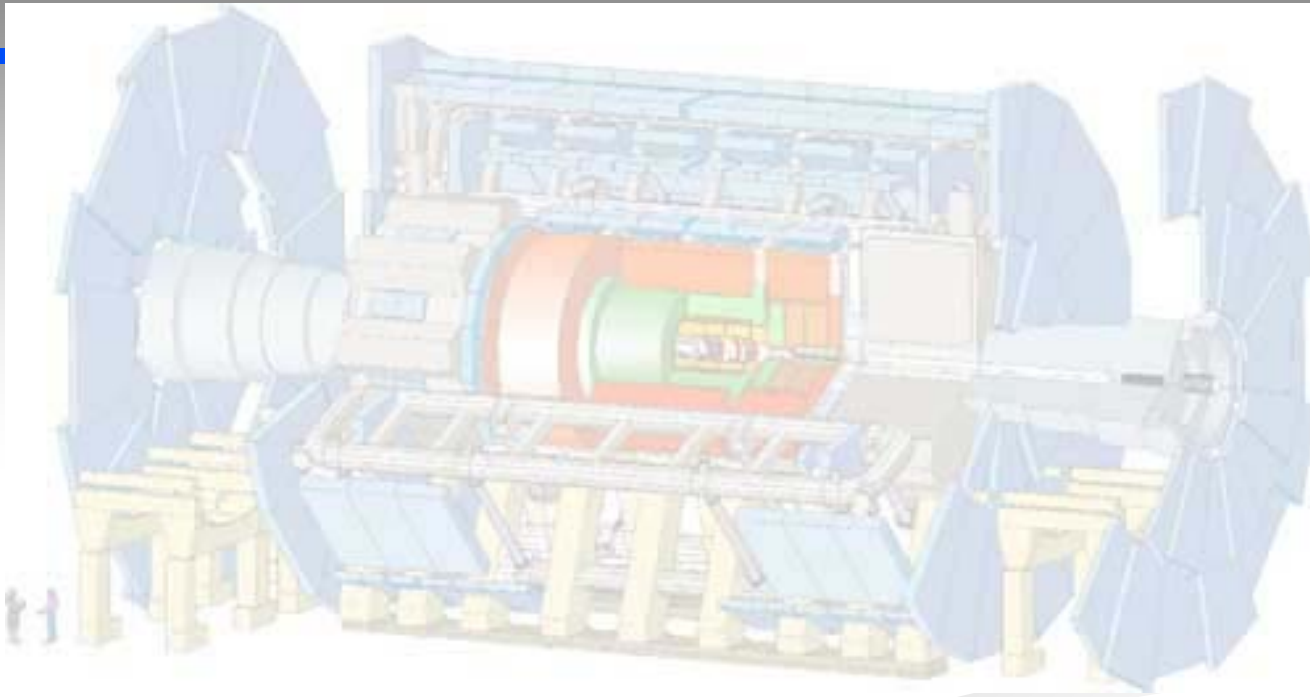
- Multi-purpose detector for LHC
- 22 m diameter, 7000 tons

- ~1850 People
- 149 Institutions, 34 Countries
- 37 Funding Agencies

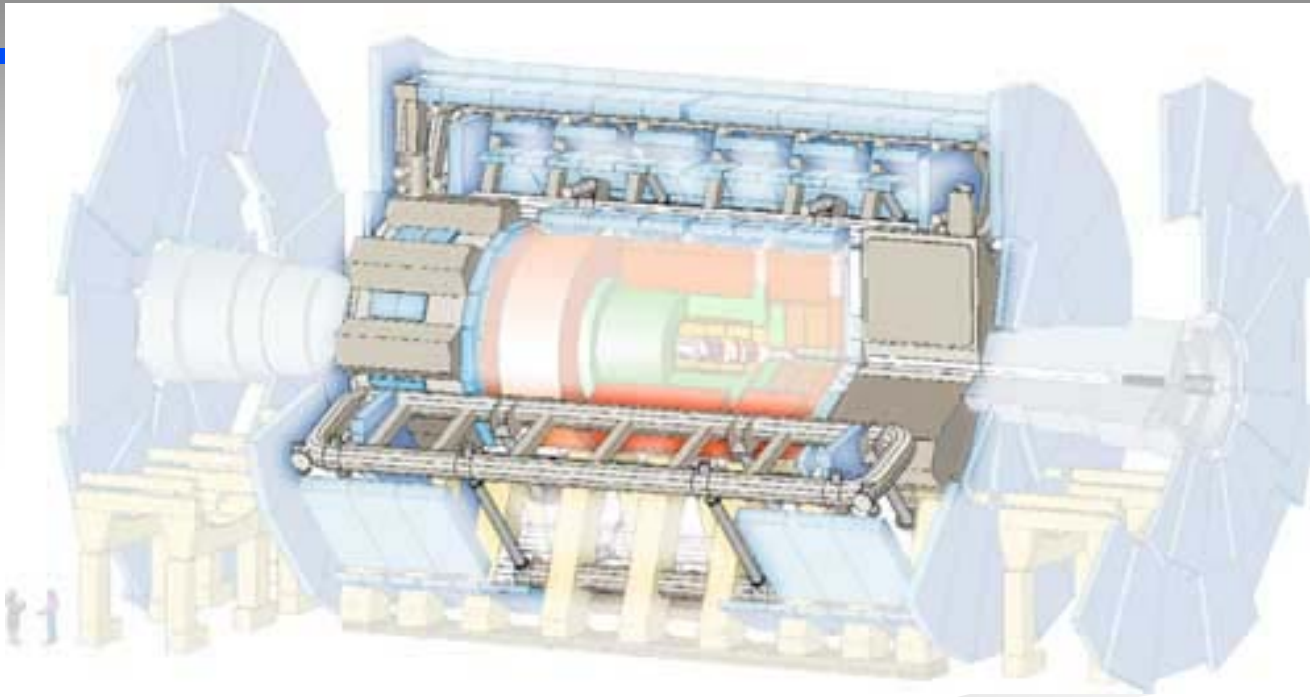


- Multi-purpose detector for LHC ➤ ~1850 People
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The ATLAS Detector



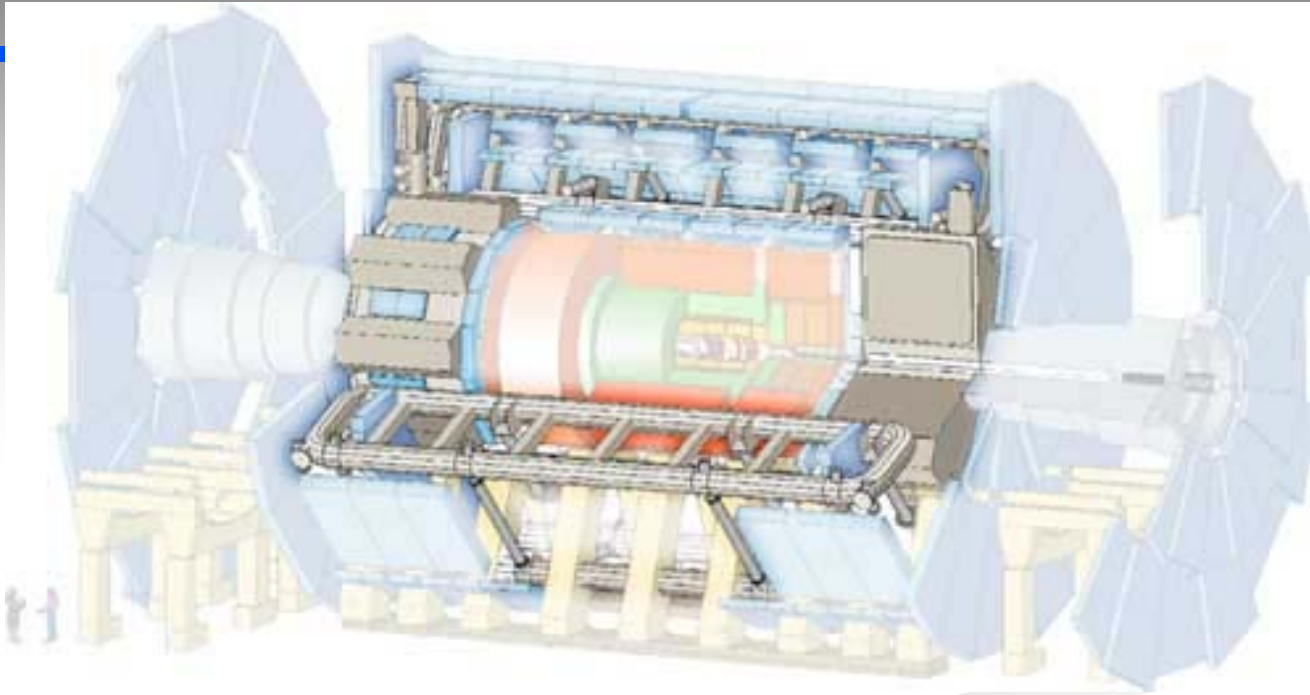
The ATLAS Detector



Magnet System

- 2T Solenoid surrounds inner detector (no field at calorimeters)
- 3.9 - 4T air core toroids for muon system

The ATLAS Detector

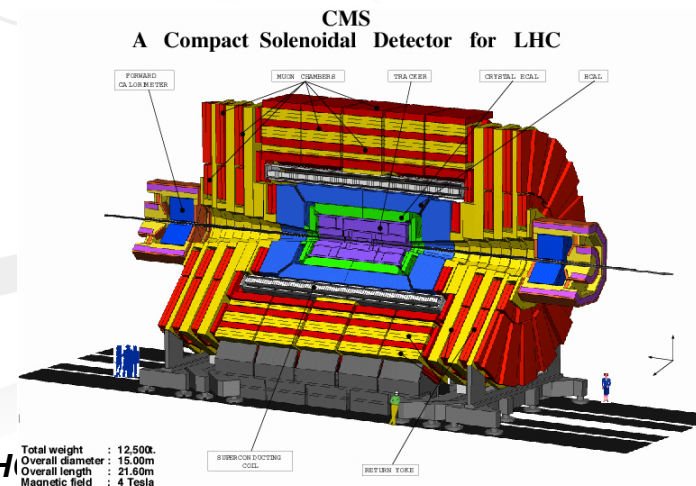


Magnet System

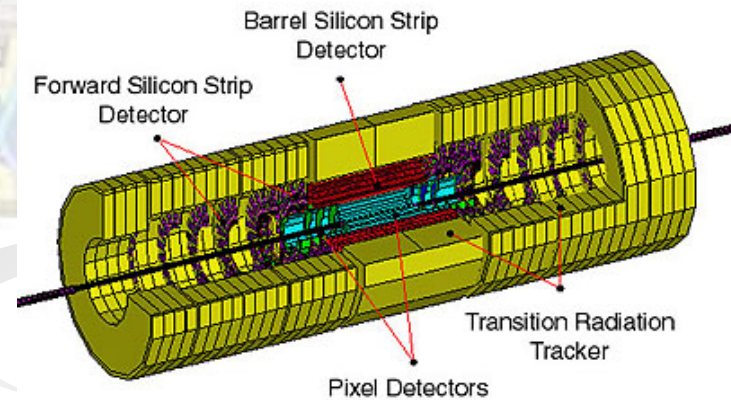
... provides names for the LHC expts

AToroidal **L**H**C** **A**pparatu**S**

VS. **C**ompact **M**uon **S**olenoid



The ATLAS Detector



Inner Tracker

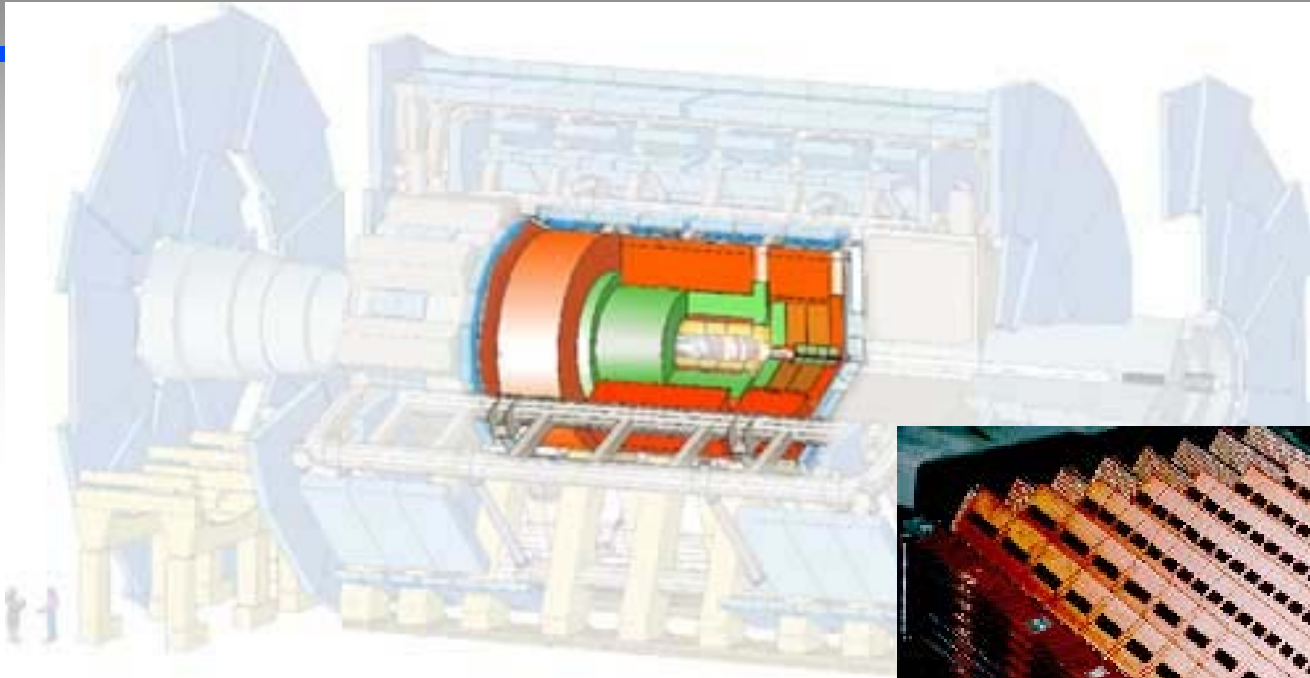
$$\frac{\sigma}{P_T} = \frac{P_T(\text{GeV})}{2000} \oplus 0.01$$

Inner Detector

- Silicon pixels and strips
- transition radiation tracker with e/π separation capabilities



The ATLAS Detector



- EM Calorimeters

Pb / LAr

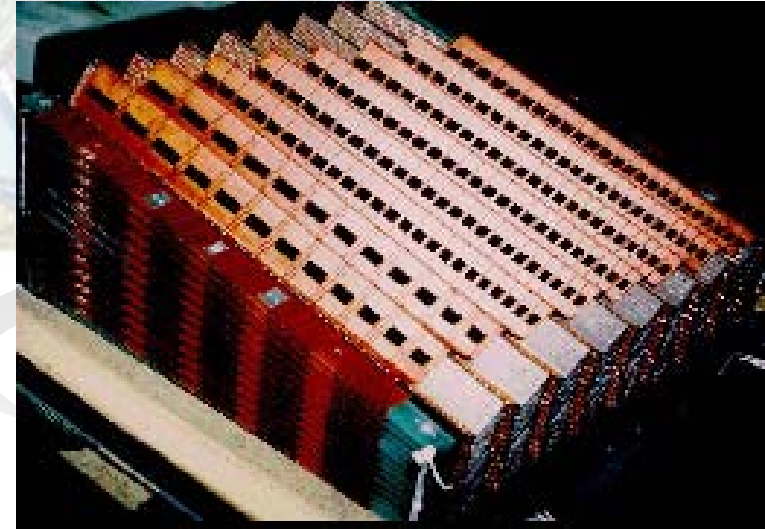
$$\frac{\sigma}{E} = \frac{10\%}{\sqrt{E(\text{GeV})}}$$

- Hadron Calorimeters

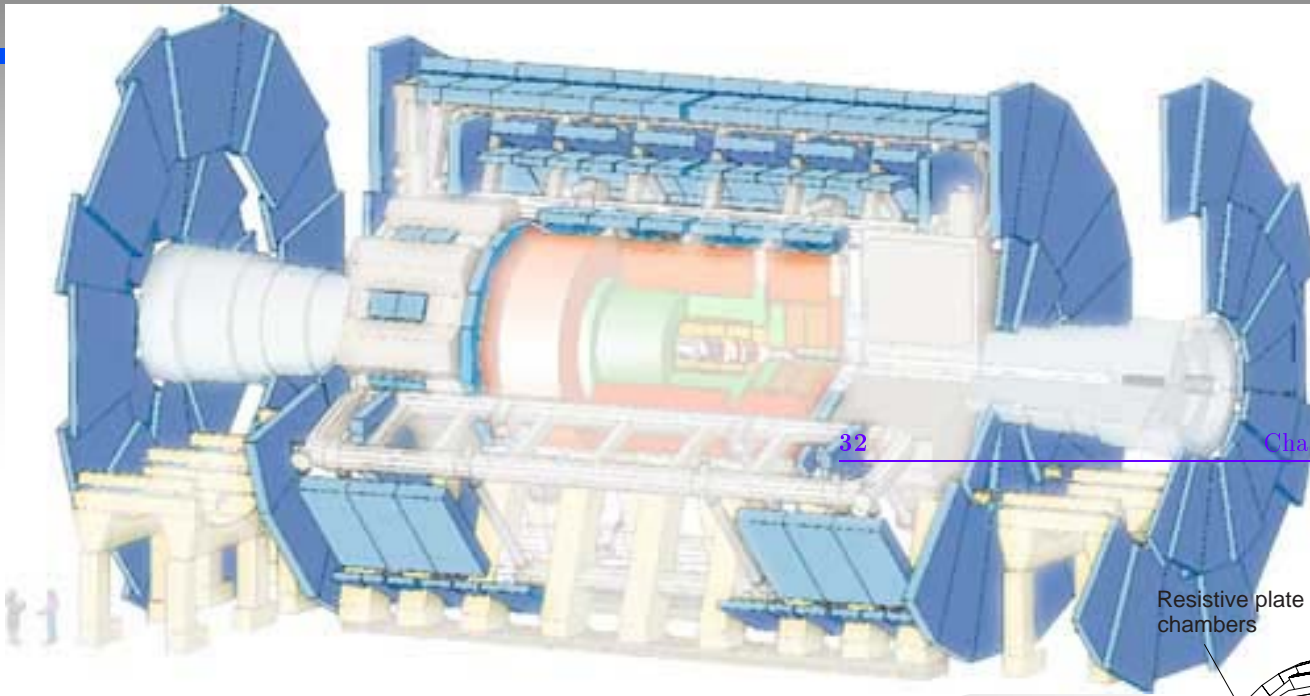
Barrel: Fe / Scintillating Tiles

Endcaps: Cu & W / LAr

$$\frac{\sigma}{E} = \frac{50\%}{\sqrt{E(\text{GeV})}} \oplus 0.03$$



The ATLAS Detector



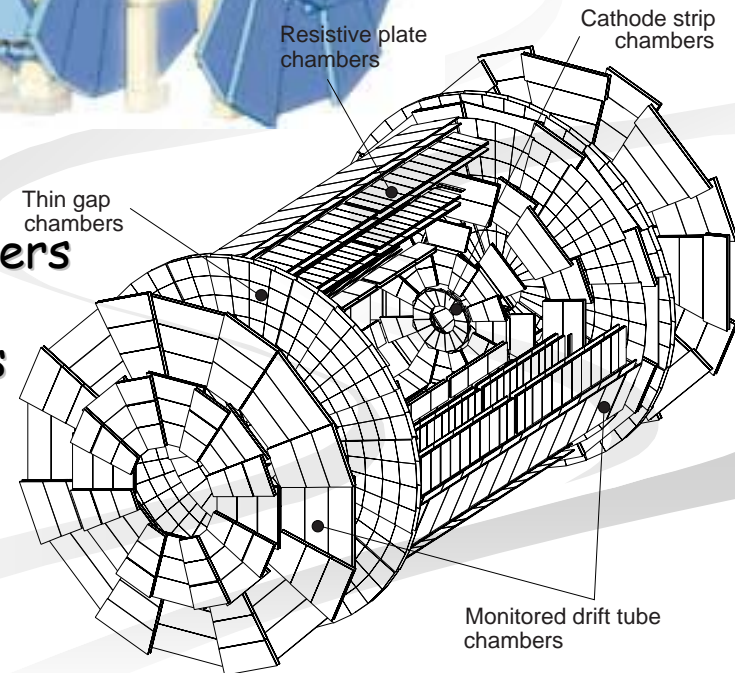
32

Chapter 3. The ATLAS Experiment

(Air Core) Muon Spectrometer

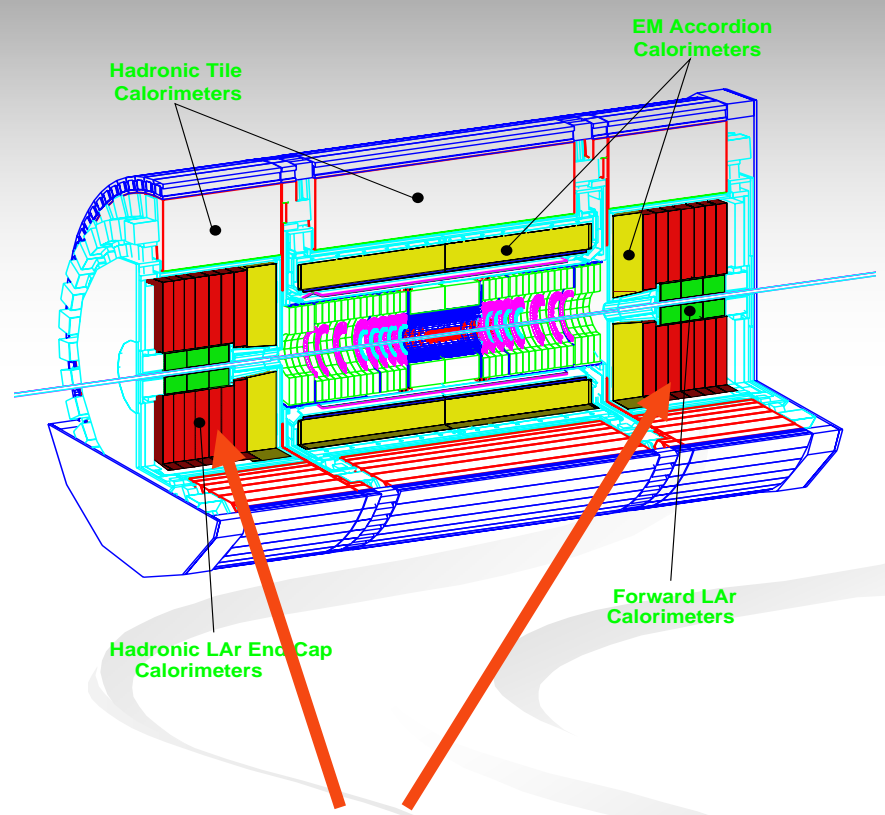
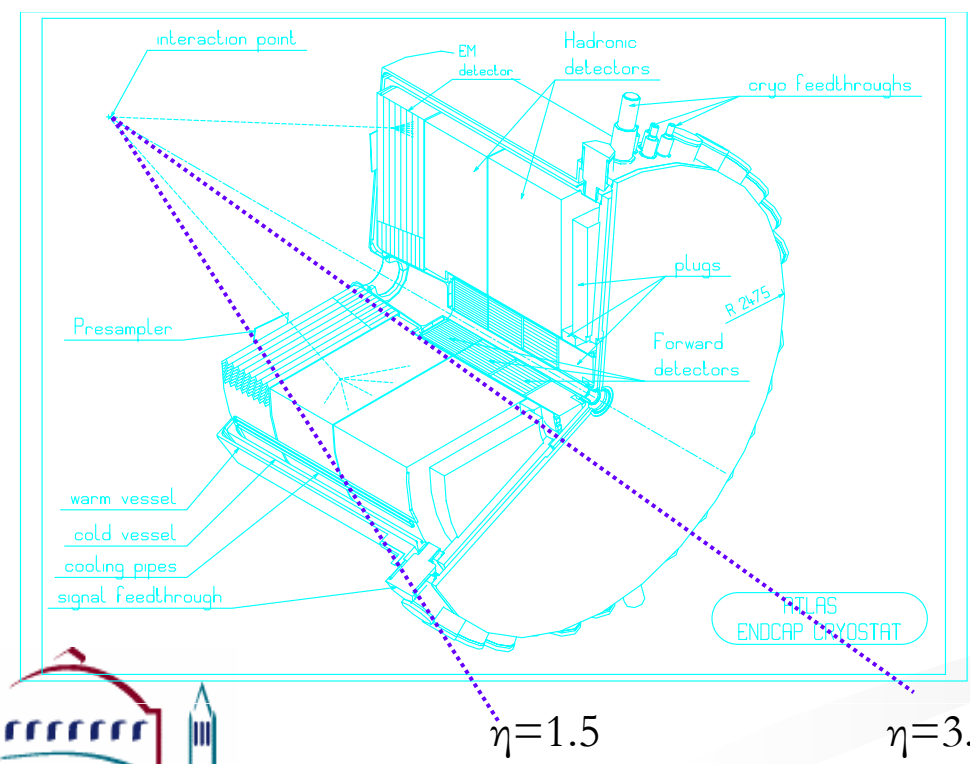
- monitored drift tubes and cathode strip chambers (precision tracking)
- resistive plate chambers and thin gap chambers (fast triggering)
- Good standalone performance

$$\frac{\sigma}{P_T} \cong 2-3\% \text{ for } P_T < 1 \text{ TeV}$$

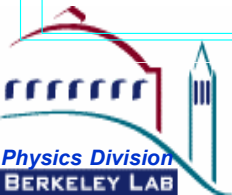




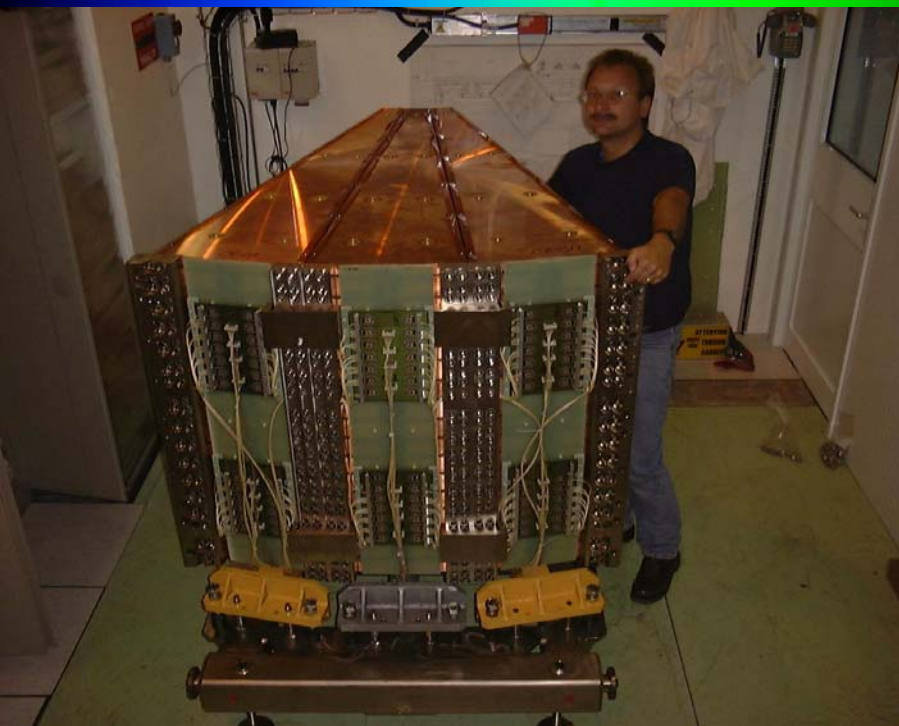
- Subdetector collaborations are busy calibrating, evaluating and understanding their detectors in beam tests.



Example: Hadronic Endcap Calorimeter



HEC Beam Test



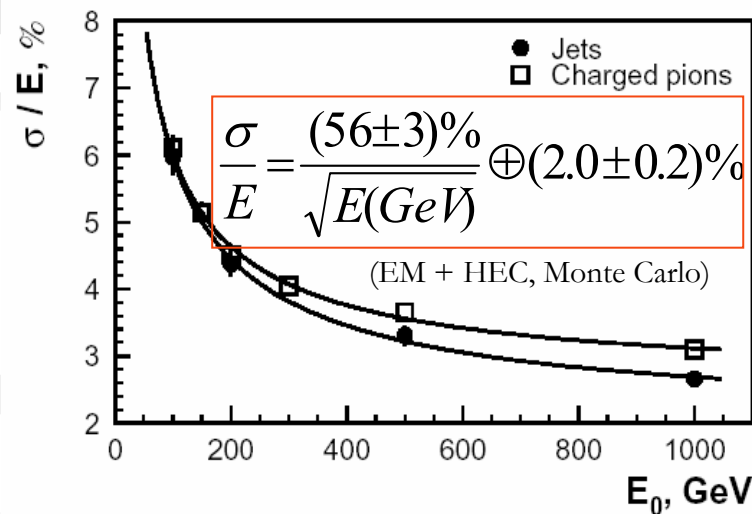
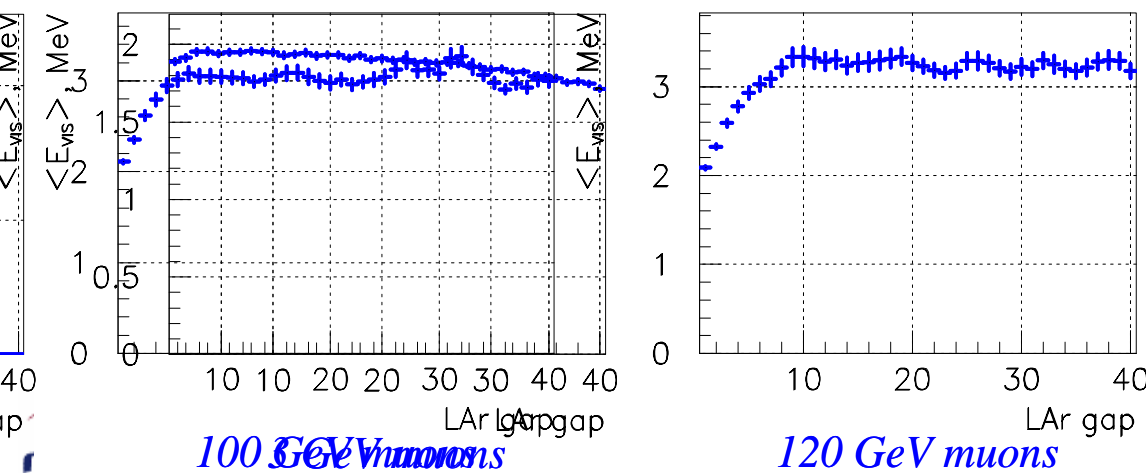
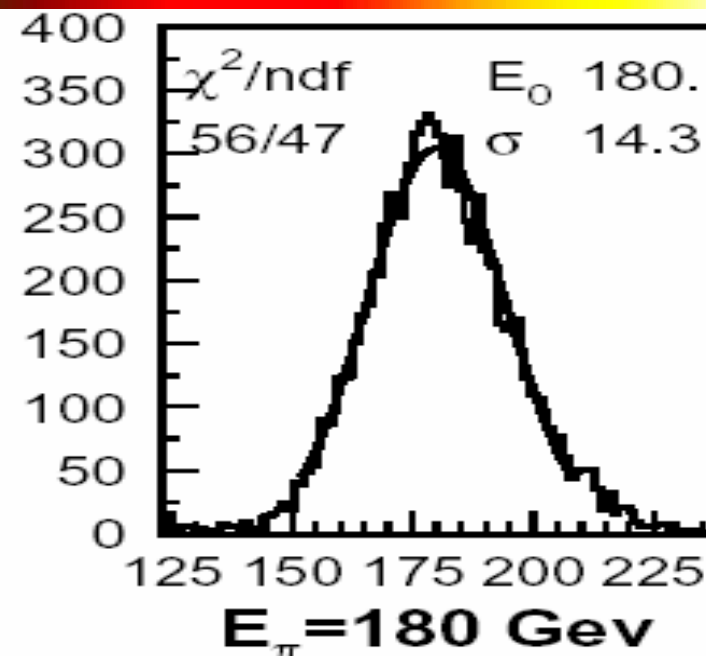
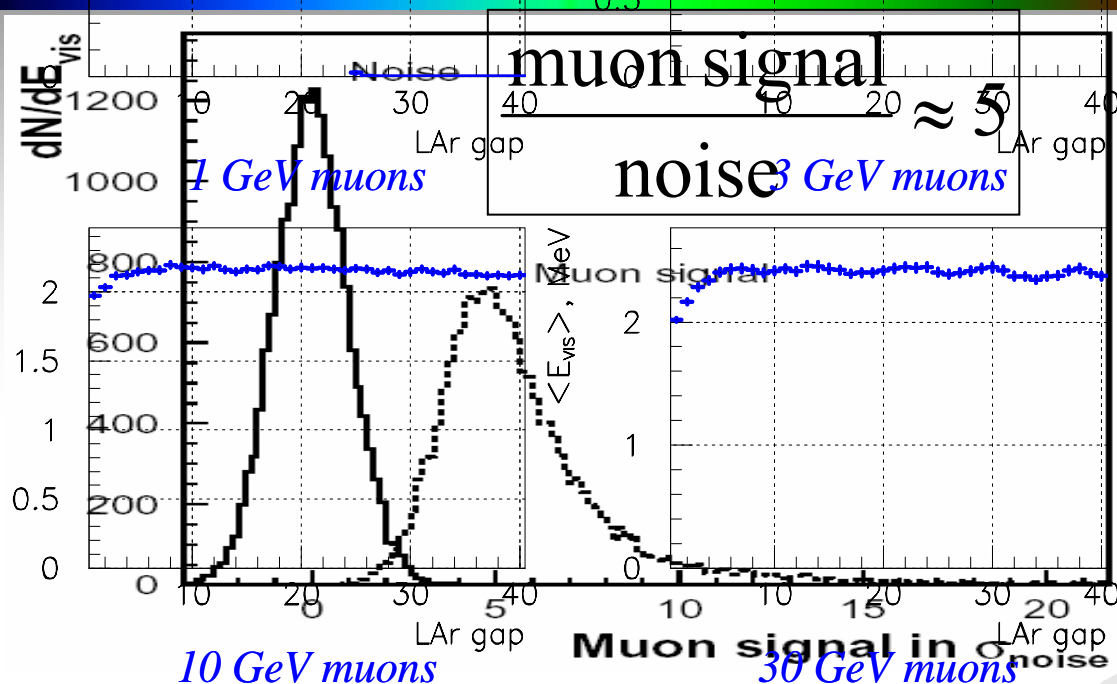
1 wheel constructed and assembled
at Triumf, Vancouver
1 wheel constructed at Dubna (Russia)
and assembled at MPI, Munich

Tested in CERN SPS $e/\mu/\pi$ beams,
1998-2001



Former USSR bomb cases

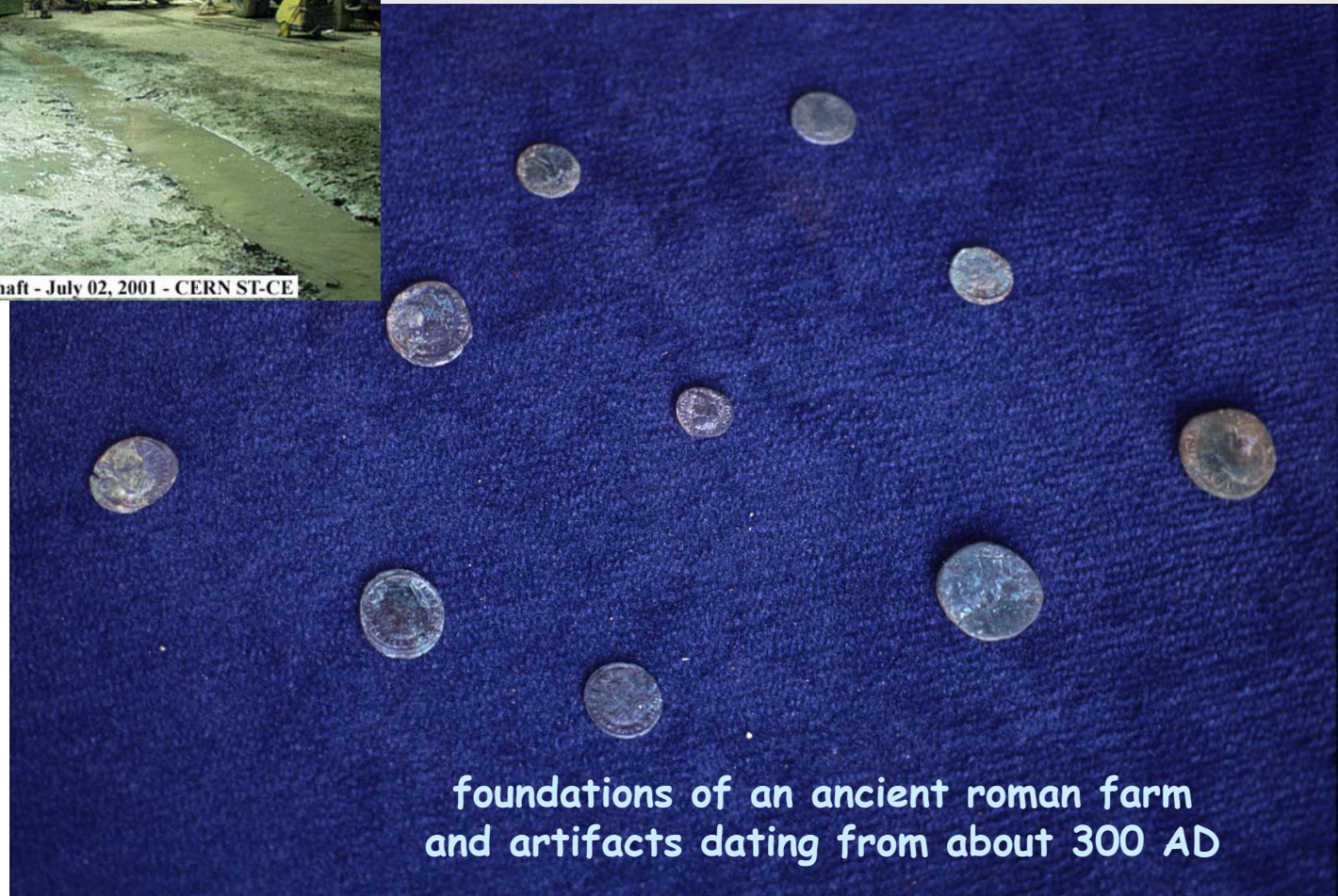
Dynamic range from MIP to few TeV



announcing the first LHC discovery



Point 1 - UX15 cavern vault; intersection with PX14 shaft - July 02, 2001 - CERN ST-CE



foundations of an ancient roman farm
and artifacts dating from about 300 AD

Gauge-boson Physics



■ Drell-Yan lepton pair production

(2 Nobel prizes... so far, 30 years of Drell Yan measurements)

$L \rightarrow$ channel for large extra Dimensions, Z'

- ✿ probe proton structure: parton density functions at small x
- ✿ Calibrate the detector

- EM energy & momentum scale from $pp \rightarrow Z^0 \rightarrow l^+l^-$

- jet energy scale from $pp \rightarrow Z^0 + \text{jet} \rightarrow l^+l^- + \text{jet}$

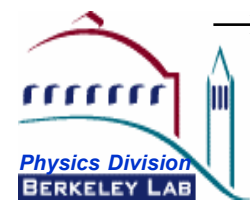
- luminosity from $pp \rightarrow Z^0, W^\pm$ event rate

→ **Important!**, our knowledge of this process feeds into the systematic errors for all physics measurements and searches

■ Key Precision Measurements of fundamental SM parameters

$\sin^2\theta_W$, $\text{Mass}(W)$

→ let's explore the $\sin^2\theta_W$ example...



Measuring $\sin^2\theta_W$ with A_{FB}



- $pp \rightarrow l^+l^-$ di-lepton signature is (almost) background free
- asymmetry arises from interference between neutral currents
$$\frac{d\sigma}{d\Omega} \propto \left| \gamma^* + Z^0 + (\text{New Physics!}) \right|^2$$
- constrains M_{HIGGS} and checks model consistency

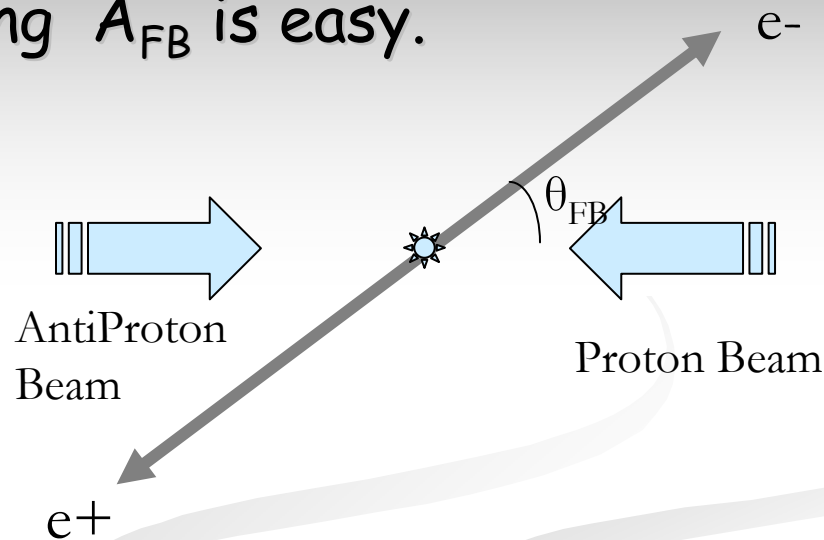
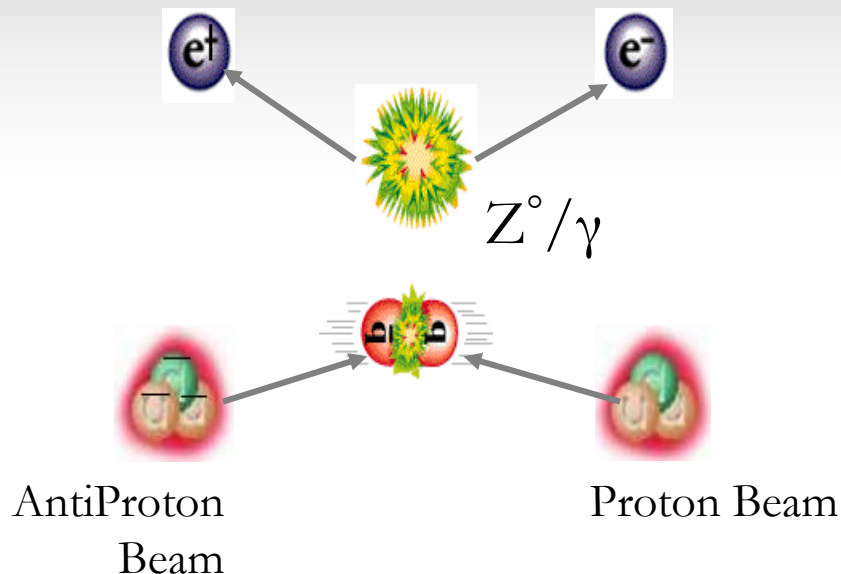
$$A_{FB} = \frac{\sigma_F - \sigma_B}{\sigma_F + \sigma_B} = b \left(a - \sin^2 \theta_{eff}^{lept}(M_Z) \right)$$

known to NLO in EW, QCD
(effects can be as large as 30%)

Measuring $\sin^2\theta_W$ with A_{FB}



- At the $\bar{p}p$ Tevatron, defining A_{FB} is easy.

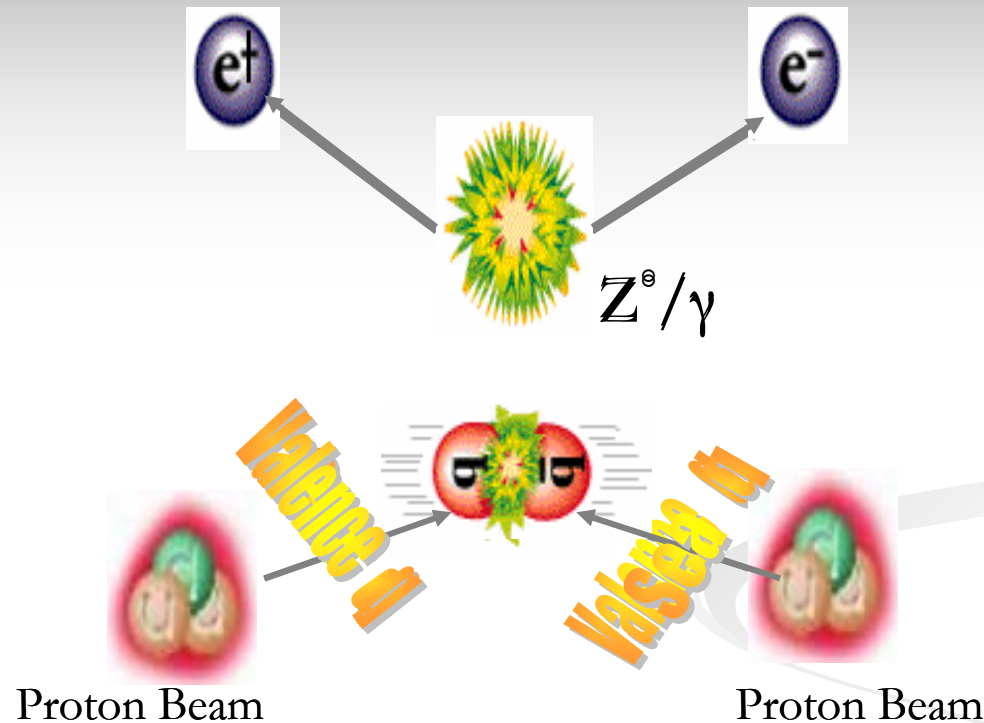


- but for symmetric proton-proton beams (LHC), there is no asymmetry WRT the beams.

Measuring $\sin^2\theta_W$ with A_{FB}



- instead, we “sign” the forward direction by the l^+l^- boost.



- measure asymmetry in charged lepton direction WRT CMS boost direction
- Asymmetry increases at high \sqrt{s}

Measuring $\sin^2\theta_W$ with A_{FB}



- Statistical precision using 100 fb^{-1} , near Z-pole ($\pm 6 \text{ GeV}$)


Cuts	A_{FB} (%)	ΔA_{FB} (%)	$\Delta \sin^2\theta_{\text{eff}}(M_Z)$
Both e^\pm , $ \eta < 2.5$	0.774	0.020	0.00066
One e^\pm , $ \eta < 2.5$ other e^\pm , $ \eta < 4.9$	1.98	0.018	0.00014

ATL-PHYS-2000-018

Sliwa, Riley, Baur

for comparison, $\Delta \sin^2\theta_{\text{eff}} = 0.00053$ combining 4 LEP expts and e, μ, τ channels [CERN-EP/2001-098]

- Performance issue: increasing forward lepton tagging acceptance greatly improves measurement
- systematic PDF uncertainty is most challenging.



Modeling our Predictions: New Monte Carlo Techniques for QCD corrections

*the real challenge for M.C. authors is
modeling subtle Standard Model effects...
new physics is (usually) easy!*

Simulating QCD Corrections

2 common approaches

Showering event Generators

(Pythia, Herwig, Isajet)

Next-to-leading order

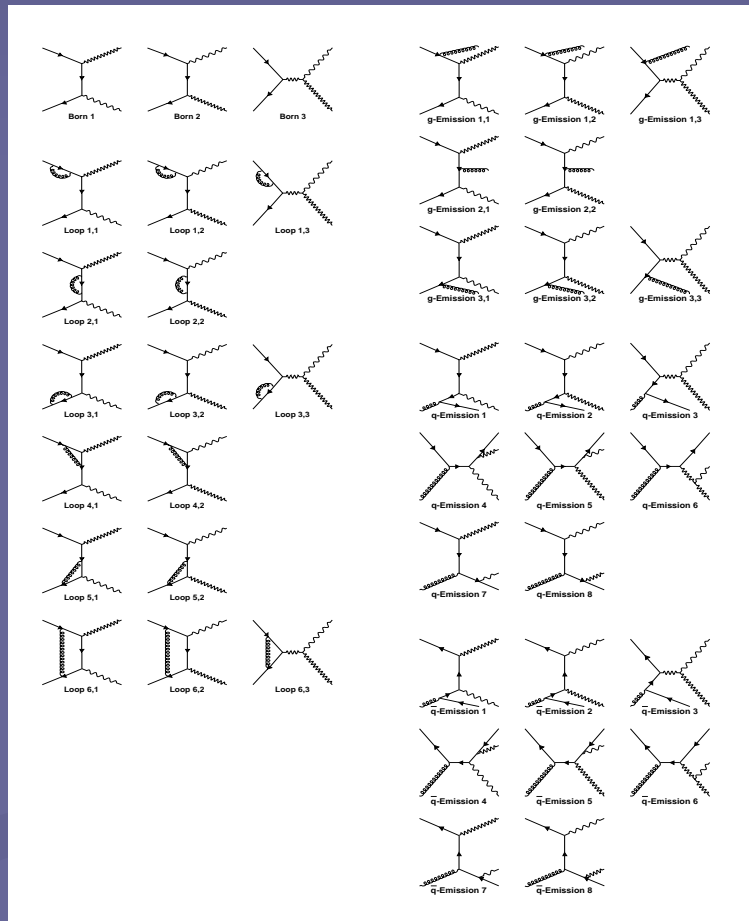
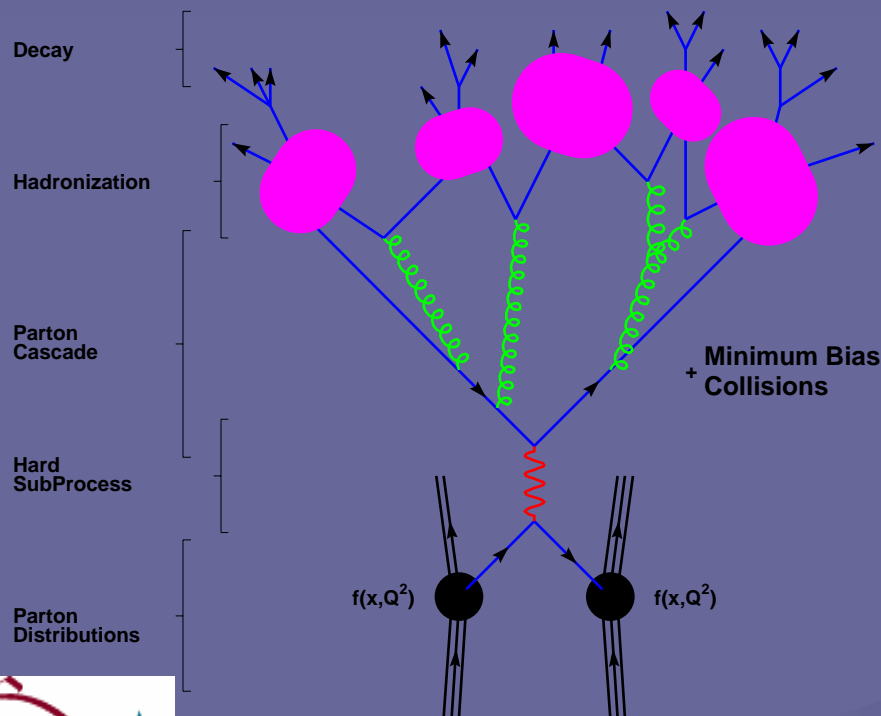
“event integrators”

Matt.Dobbs@cern.ch

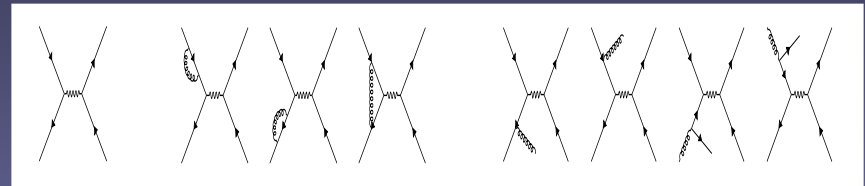
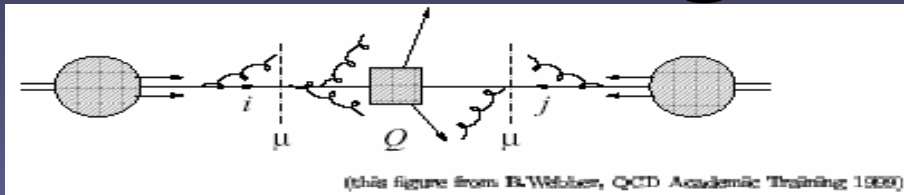
Diboson Feynman Graphs at NLO

WTDX-ed on June 18, 2001

3



Simulating QCD Corrections



Showering Event Generators

- ☺ Event generation is probabilistic... freq predicted by theory.
- ☺ exclusive prediction
 - you get the whole event record
- ☺ all orders approximation of multiple emissions
- ☺ valid in soft/collinear emission regions
- ☹ not valid for hard, well separated partons
- ☹ Normalization is LO

NLO Matrix Elements

- ☺ good prediction of hard central emissions
- ☺ best prediction of total σ
- ☹ one order in α_s
 - at most one "jet"
- ☹ fixed order perturbation is not valid for small $P_T(\text{jet})$
- ☹ event weights are negative (unphysical) in some phase space regions

Phase Space Veto Method

Implemented as an event generator for

$$p\bar{p} \rightarrow Z^0 + X \rightarrow l^+ l^- + X$$

but everything *applies in general to any colour singlet production* process at hadron colliders



(Matt: Click here if you are short on time!)

NLO 'event integrators' $pp \rightarrow Z+X \rightarrow l^+l^-+X$

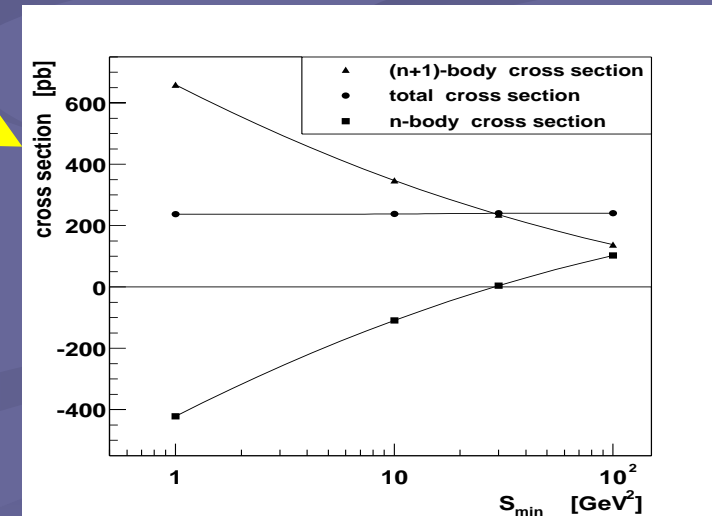
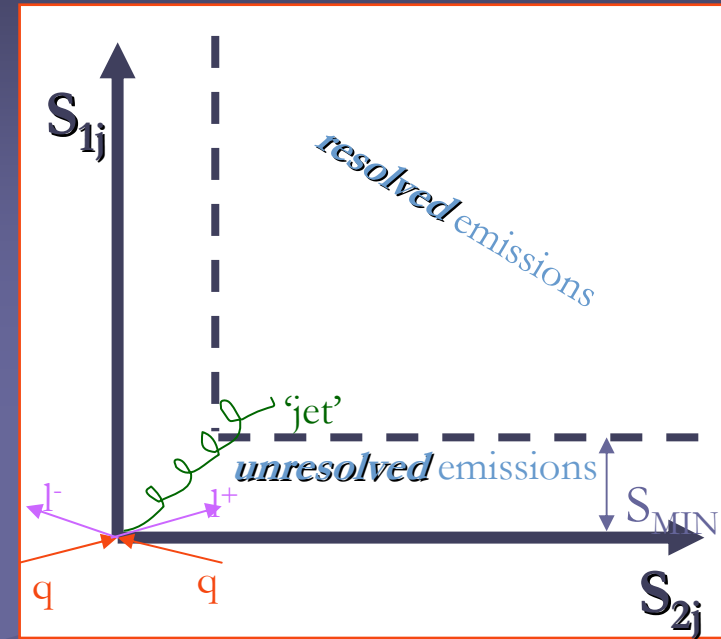
Regularization scheme example:

Phase Space Slicing ("S_{MIN} slicing")

- partition phase space:
 - **resolved** region: integrated numerically
 - **unresolved** region: integrated analytically
- programmed as two separate generators
- cross section is independent of our S_{MIN} choice

$$\sigma^n(S_{MIN}) + \int_{S_{ik} > S_{MIN}} \sigma^{n+1}(\Phi_{+1}) d\Phi_{+1} = \text{Const}$$

→ can choose (almost) any S_{MIN} we like.



Phase Space Veto Method

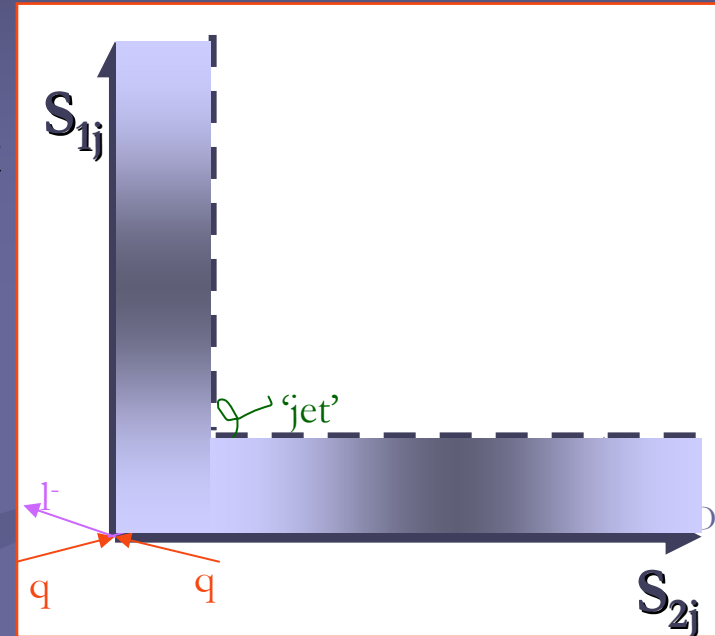
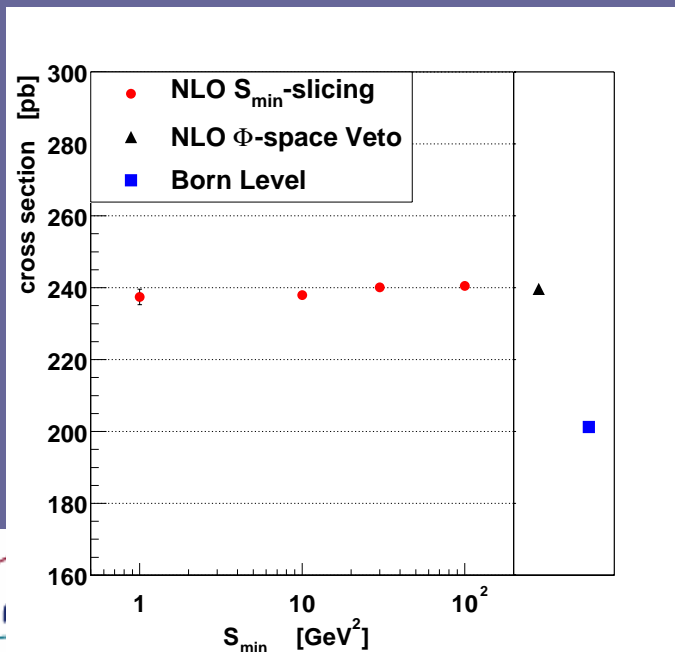
Dobbs, Phys.Rev.D64,034016 (2001), Phys Rev D65,094011 (2002)

Pötter, Phys.Rev.D 63,114017 (2001) [DIS]

Recall: \rightarrow can choose (almost) any S_{MIN} we like

$$\sigma^n(S_{\text{MIN}}) + \int_{S_{ik} > S_{\text{MIN}}} \sigma^{n+1}(\Phi_{+1}) d\Phi_{+1} = \text{Const}$$

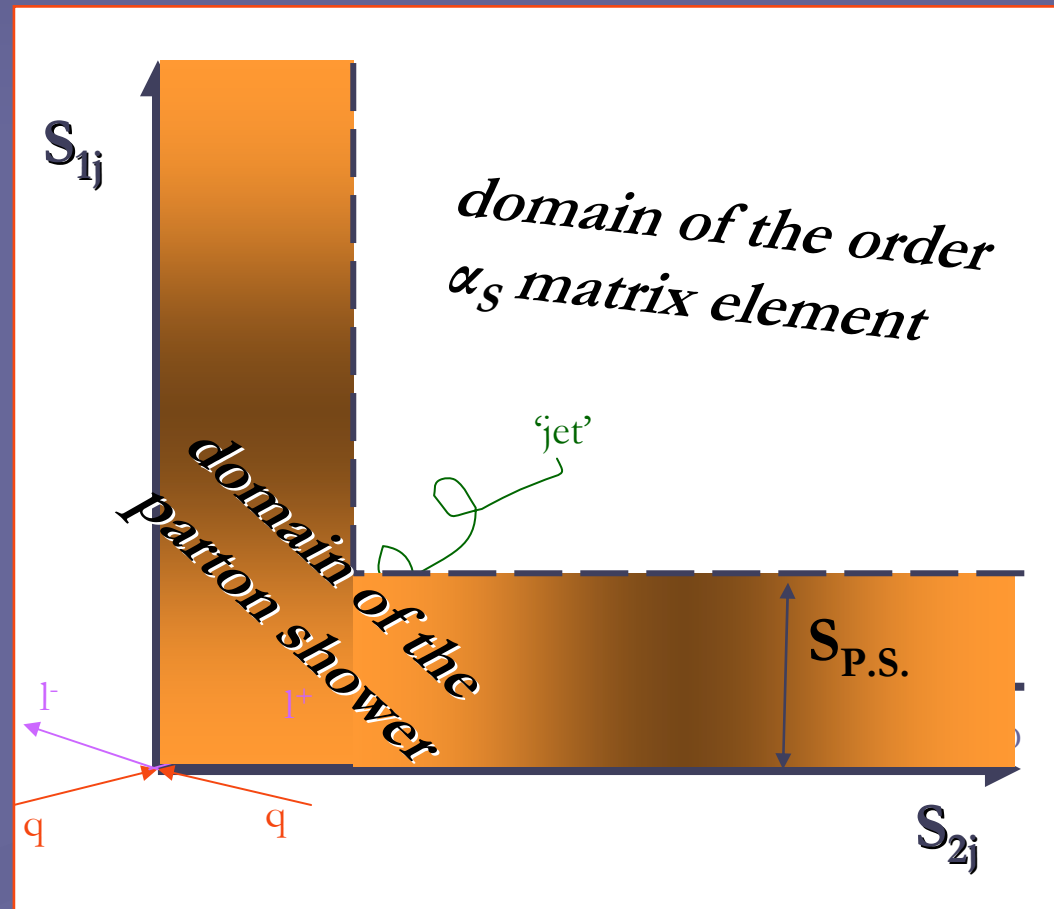
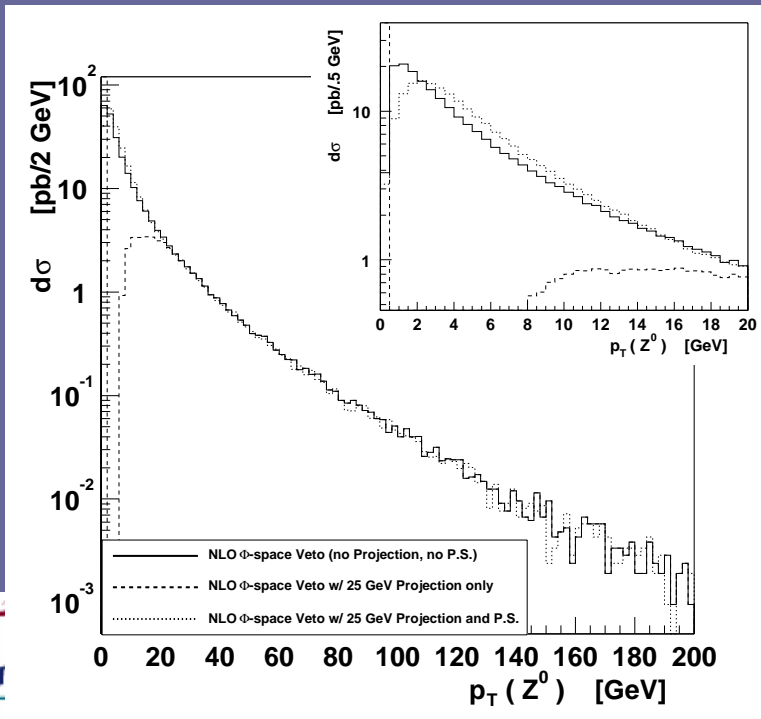
choose S_{ZERO}
i.e. unresolved contribution, $\sigma^n(S_{\text{ZERO}})=0$
on an event-by-event basis.



Addresses the first issue, because it carves out the region of negative weights \rightarrow i.e. it allows us to re-formulate the NLO calculation into a true (probabilistic) event generator.
 \rightarrow while maintaining the reduced scale dependence provided by the NLO calculation.

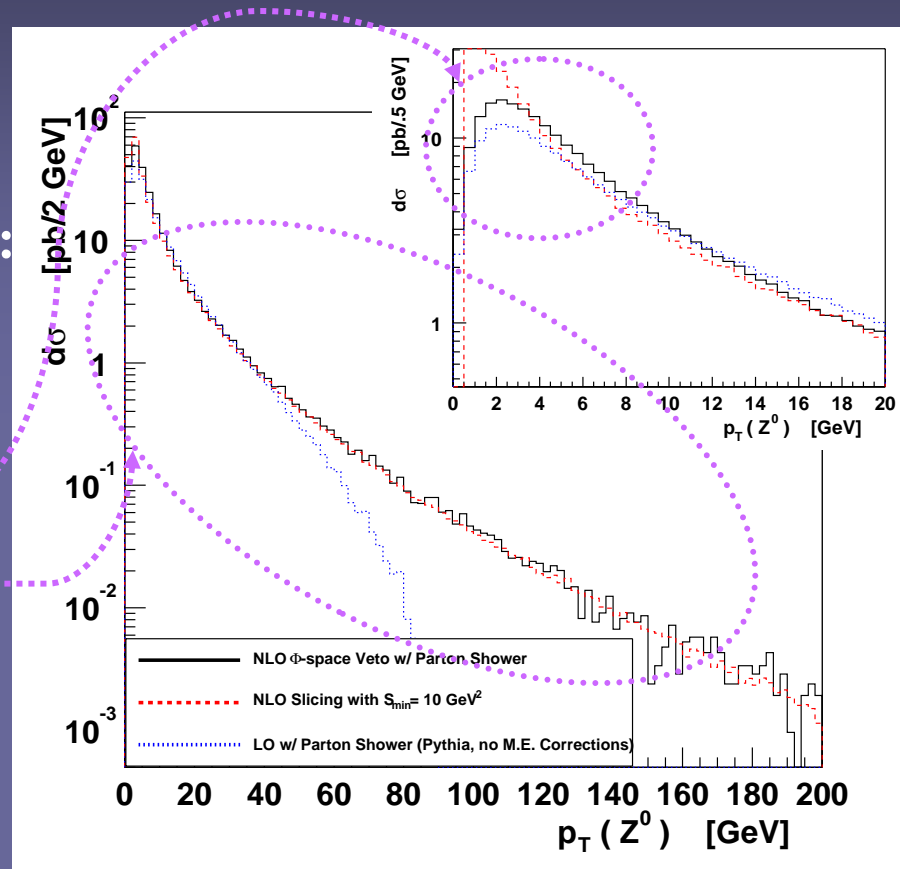
Phase Space Veto Method: shower evolution

- our description of the hard central region is dominated by the NLO matrix elements
- ideally, we want the small P_T region to be the domain of the Parton Shower



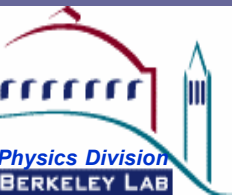
Phase Space Veto Method

- event generator in the true sense
 - you get unweighted events
 - well suited for expt. applications:
 - detector simulation,
 - hadronization, etc.
- dominated by parton shower in the soft/collinear regions
- dominated by $O(\alpha_s)$ matrix element in hard/central regions
- normalization is NLO, maintains reduced scale dependence.



Phase Space Veto Method

- attacks a phenomenological issue from an experimental viewpoint
- implemented as a new $pp \rightarrow Z^0/\gamma^* + X \rightarrow l^+l^- + X$ event generator
 - ⊗ see Dobbs, Phys Rev D65,094011 (2002)
 - ⊗ (uses LUND parton shower)
 - ⊗ efficiency and event generation time competes with L.O. generators
 - ⊗ another step towards modularized event generators (HepMC, HepUP)
 - ⊗ written in Object Oriented C++
- preliminary results indicate there will be further benefits for more complicated processes like diboson production.
[see Dobbs, Phys Rev D64,034016 (2001)]
- The phenomenology community is listening!
 - ↪ 'competing' version from Herwig announced February 2002: S.Frixione/B.Webber (hadronic diboson production at NLO)
 - ↪ Minami-Tateya (KEK) group (GRACE), Yoshimasa Kurihara, $pp \rightarrow Z + X$
 - ↪ Lund has new effort underway: S. Burby/T. Sjöstrand (ME corrections to $W\gamma$ WW WZ)

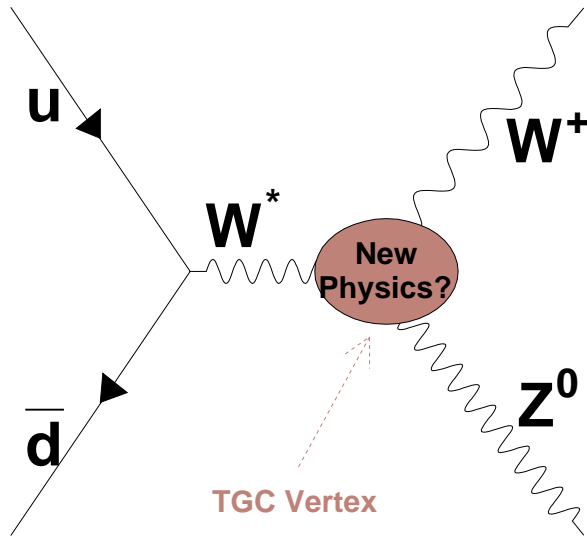


... Modeling our Predictions
New Monte Carlo Techniques
for QCD corrections



I

Probing the Triple Gauge-boson Couplings



- non-abelian $SU(2)_L \times U(1)_Y$ gauge group (foundation of SM!)
 $\rightarrow WW_\gamma \quad WWZ$ couplings

- most-general C & P conserving WWZ, WW_γ vertices are specified by just **5 parameters**:

$$\underbrace{\Delta g_Z^1, \Delta \kappa_Z, \Delta \kappa_\gamma}_{\text{grow like } \sqrt{\hat{s}}} \quad \underbrace{\lambda_Z, \lambda_\gamma}_{\text{grow like } \hat{s}} \equiv \text{ZERO in the S.M.}$$

big advantage for LHC

\rightarrow model independent parameterization

- **Probe tool**: sensitive to low energy remnants of new physics operating at a higher scale
- **complement** to direct searches

Probing the Triple Gauge-boson Couplings



- theoretical expectation for new physics at 1 TeV

anomalous TGC's at most
for new physics at 1 TeV
[hep-ph/9503425 DPF]

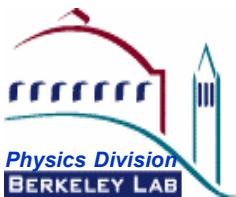
$$\mathcal{O}\left(\frac{M_W^2}{\Lambda_{N.P.}^2}\right) \approx \mathcal{O}(0.01)$$

- LEP^{Combined,ICHEP2000} & Tevatron^{Expected, RunII}

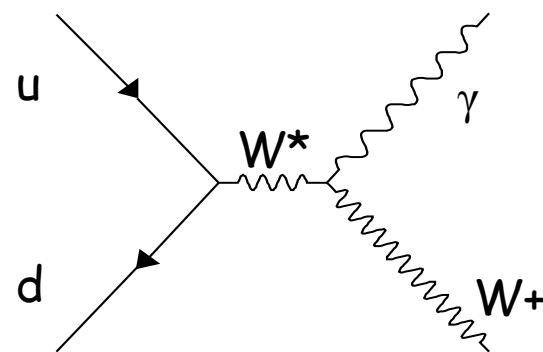
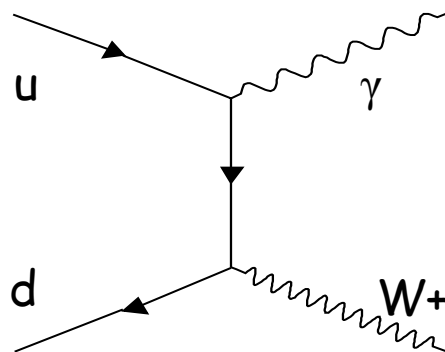
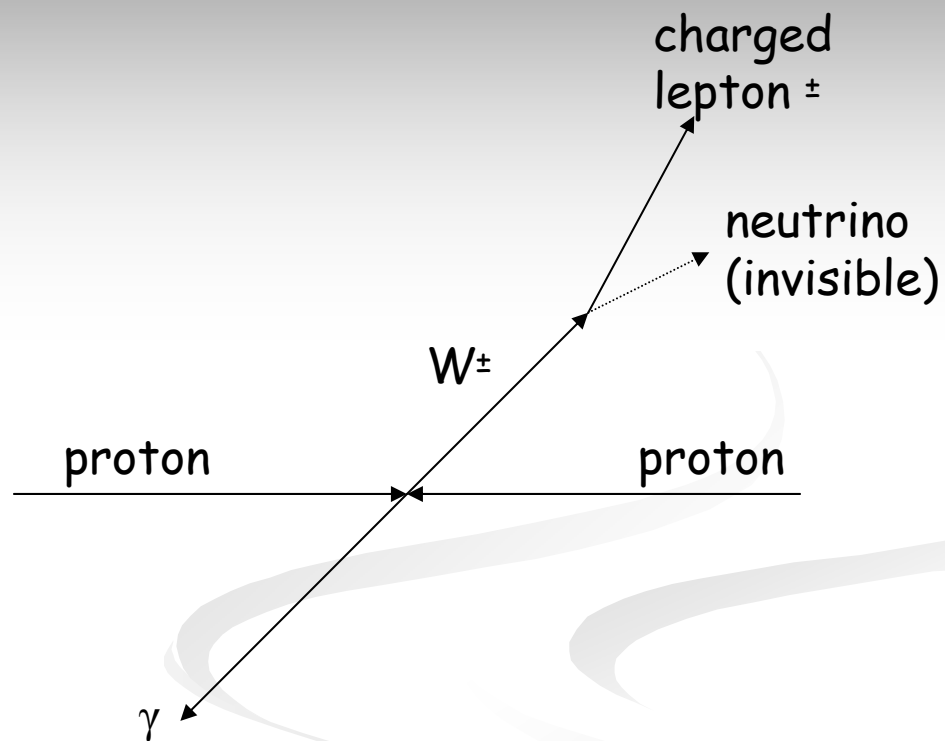
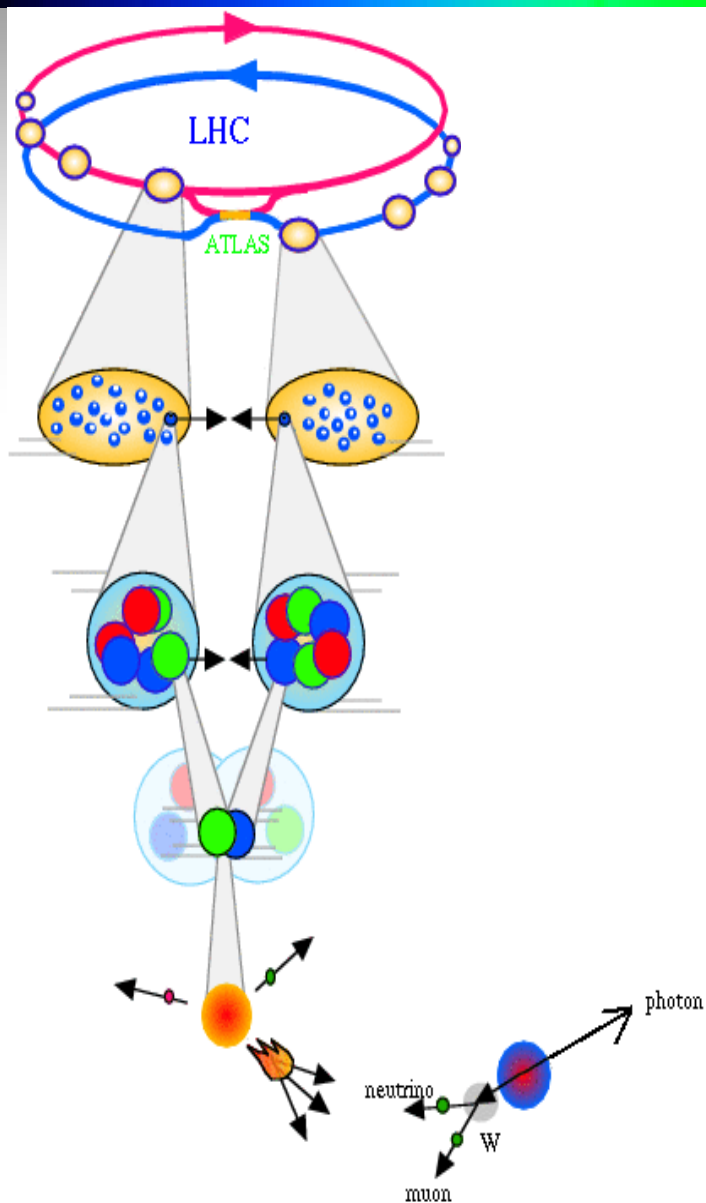
$$\mathcal{O}(.10-.20)$$

95% Confidence Limits

- probe the WWZ, WW_γ vertices with *leptonic decay channels* of WZ and W_γ production



$$pp \rightarrow W^\pm \gamma \rightarrow l^\pm \nu \gamma$$



W γ production at LHC



Consider leptonic decay channels only: $e^\pm \nu \gamma$, $\mu^\pm \nu \gamma$

Number of Events for 30 fb⁻¹

	$Z\gamma$	$W+\text{jet}$	$Z+\text{jet}$	$t\bar{t}(\gamma)$	$b\bar{b}(\gamma)$	$\gamma+\text{jet}$	$W \rightarrow l\nu\gamma$	$W\gamma \rightarrow \tau\nu\gamma$	All Backgrounds	$W\gamma$ Signal	$\frac{S}{B}$
preselection	2436	4367	7398	1561	253	956	20	710	17701	17717	1.0
$P_T^{\gamma} > 100 \text{ GeV}$	1277	2097	2101	945	160	894	14	665	8153	10638	1.30
$P_{1\pm}^{\gamma} > 25 \text{ GeV}$	1196	1938	1800	837	64	664	13	586	7098	10066	1.42
$P_{\text{miss}}^T > 25 \text{ GeV}$	377	1557	215	689	43	44	12	574	3511	7311	2.08
$\Delta R(\gamma, l^\pm) < 1$	376	1543	183	611	42	44	12	574	3385	6791	2.01
$\Sigma_{\text{jets}} P_{\text{jet},i}^T < 100 \text{ GeV}$	341	1280	133	286	26	11	12	534	2623	4262	1.62

jets faking photons is largest background
→ highlights particle ID performance

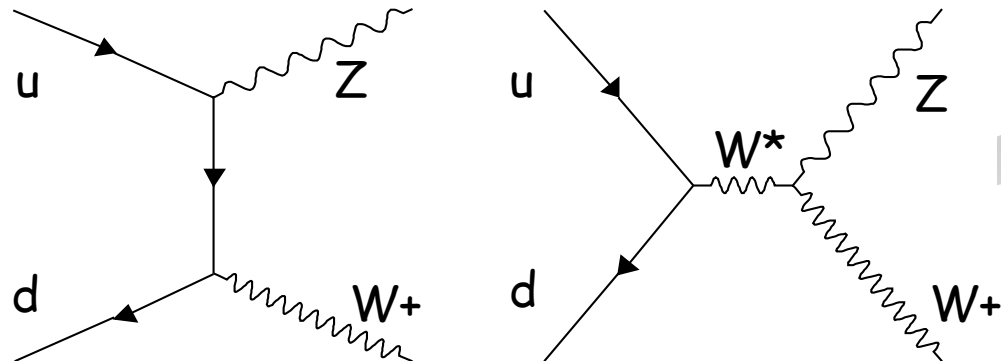
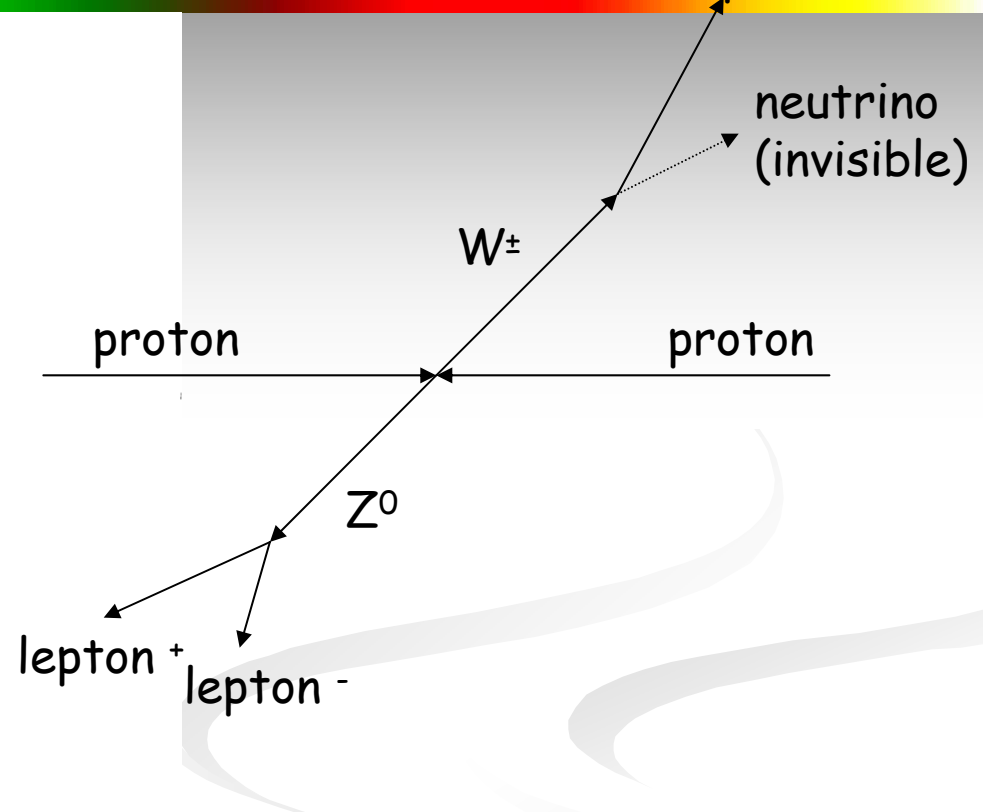
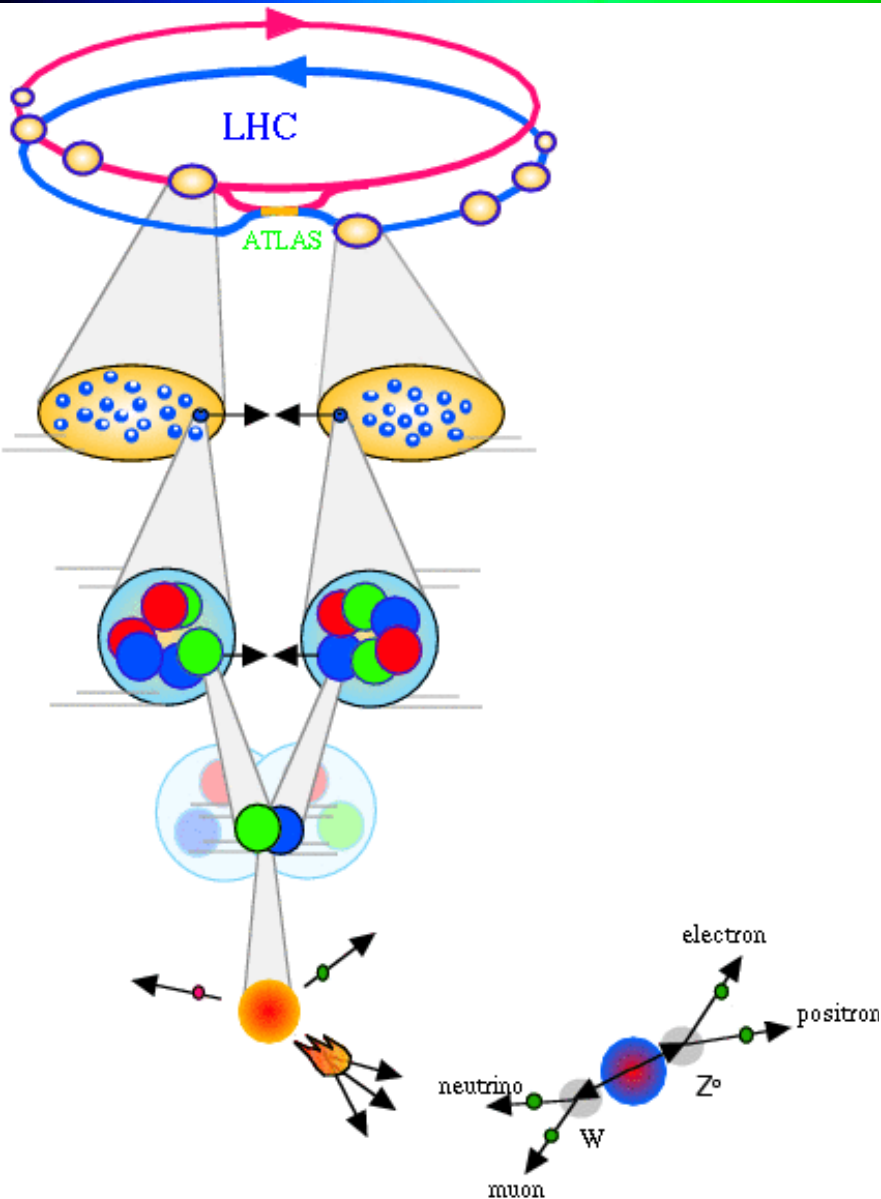
unlike the Tevatron,
 $W \rightarrow \tau \nu$ is significant
background for LHC

- Backgrounds: LO \times (k=1.5)
- Signal: NLO(α_s)

- ⊗ cuts designed for purity are optimized at LO
- cuts which isolate the phase space where TGC diagrams dominate (i.e. address $\mathcal{O}(\alpha_s)$ effects) are chosen so as to optimize the confidence limits.

$$pp \rightarrow W^{\pm} Z^0 \rightarrow l^{\pm} \nu l^{+} l^{-}$$

charged
lepton $^{\pm}$



Backgrounds to WZ production



Number of Events for 30 fb⁻¹

	# events			All Backgrounds	WZ Signal	$\frac{S}{B}$	Spread in Stat. 95% C.L.		
	Z+jet	ZZ	$t\bar{t}$				λ_Z	$\Delta\kappa_Z$	Δg_Z^1
preselection	631	576	745	1952	3663	1.88	0.014	0.29	0.020
3 leptons, $P_{1\pm}^T > 25$ GeV	398	500	461	1359	3285	2.42	0.014	0.29	0.020
$P_{\text{miss}}^T > 25$ GeV	3.2	90	357	450	2453	5.44	0.014	0.28	0.019
$ M_{l+l-} - M_Z < 10$ GeV	2.8	76	65	144	2331	16.2	0.014	0.29	0.020
$\Sigma_{\text{jets}} P_{\text{jet}_i}^T < 100$ GeV	2.5	72	44	119	1987	16.7	0.013	0.23	0.016

almost background free

statistical limits depend very weakly on obtaining good purity.

- Backgrounds: LO \times (k=1.5)
- Signal: NLO(α_s)

⌘ bb and Z γ backgrounds are negligible

Extracting the confidence intervals



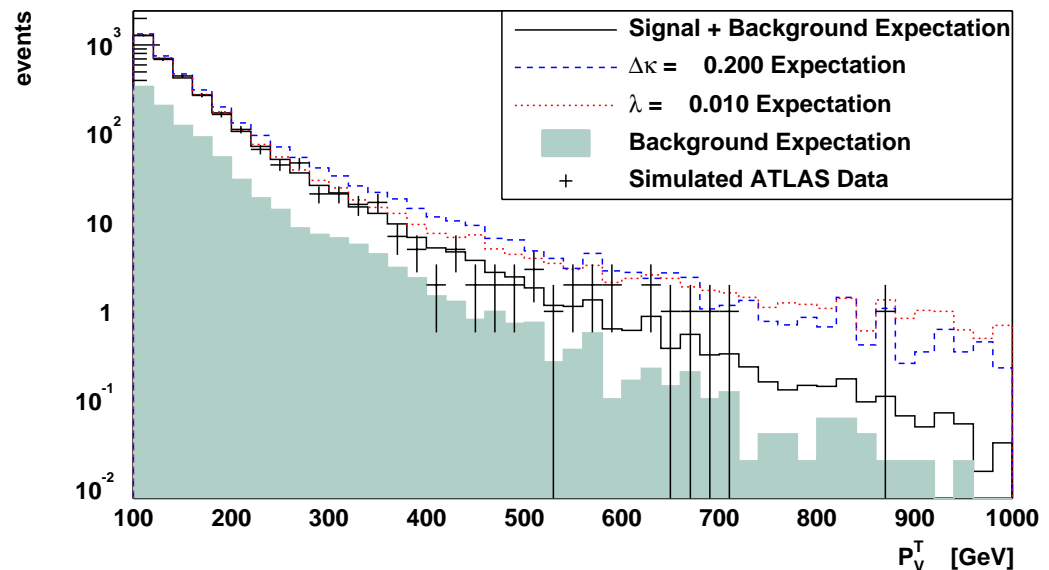
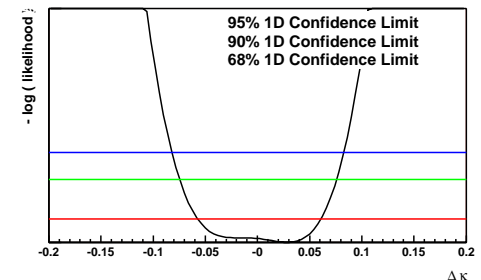
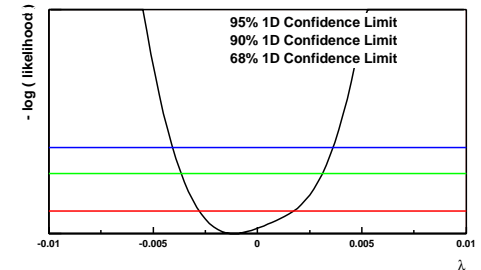
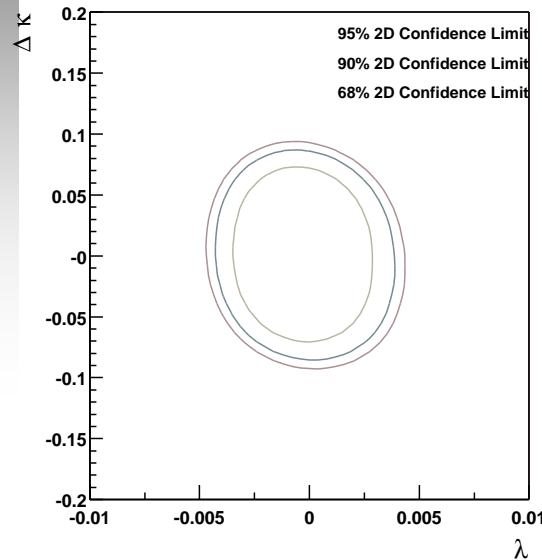
- binned max. likelihood fit to $P_T(V)$ distribution

- ⊗ robust → no need to reconstruct CMS vectors
- ⊗ directly measured

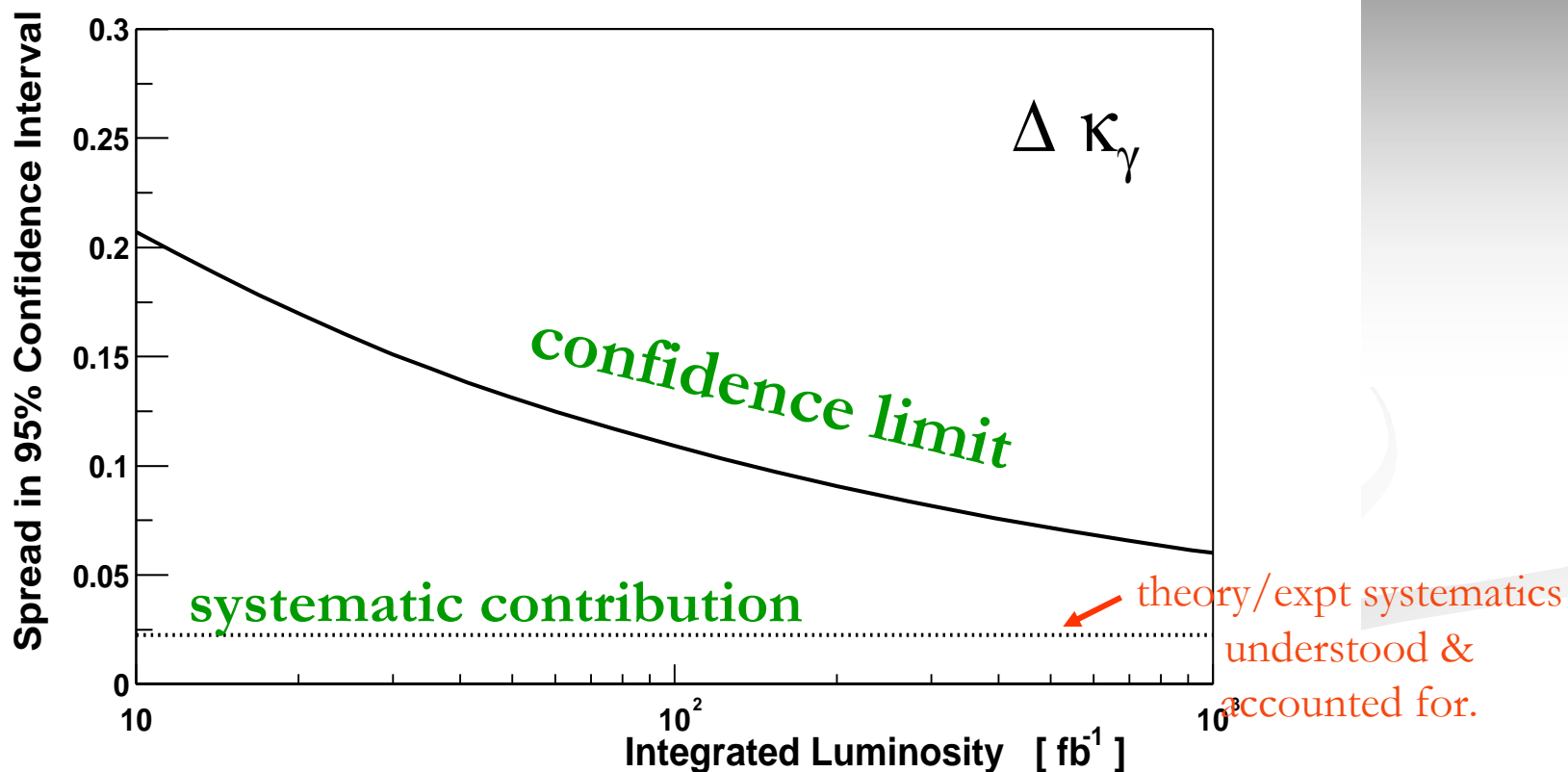
- investigated :

- ⊗ optimal observables
- ⊗ multi-variant fits [$P_T(V) \times P_T(A_W)$ is best]
- ⊗ other 1-D distributions

- luminosity systematic avoided by considering distribution shapes only



Limits vs. Integrated Luminosity

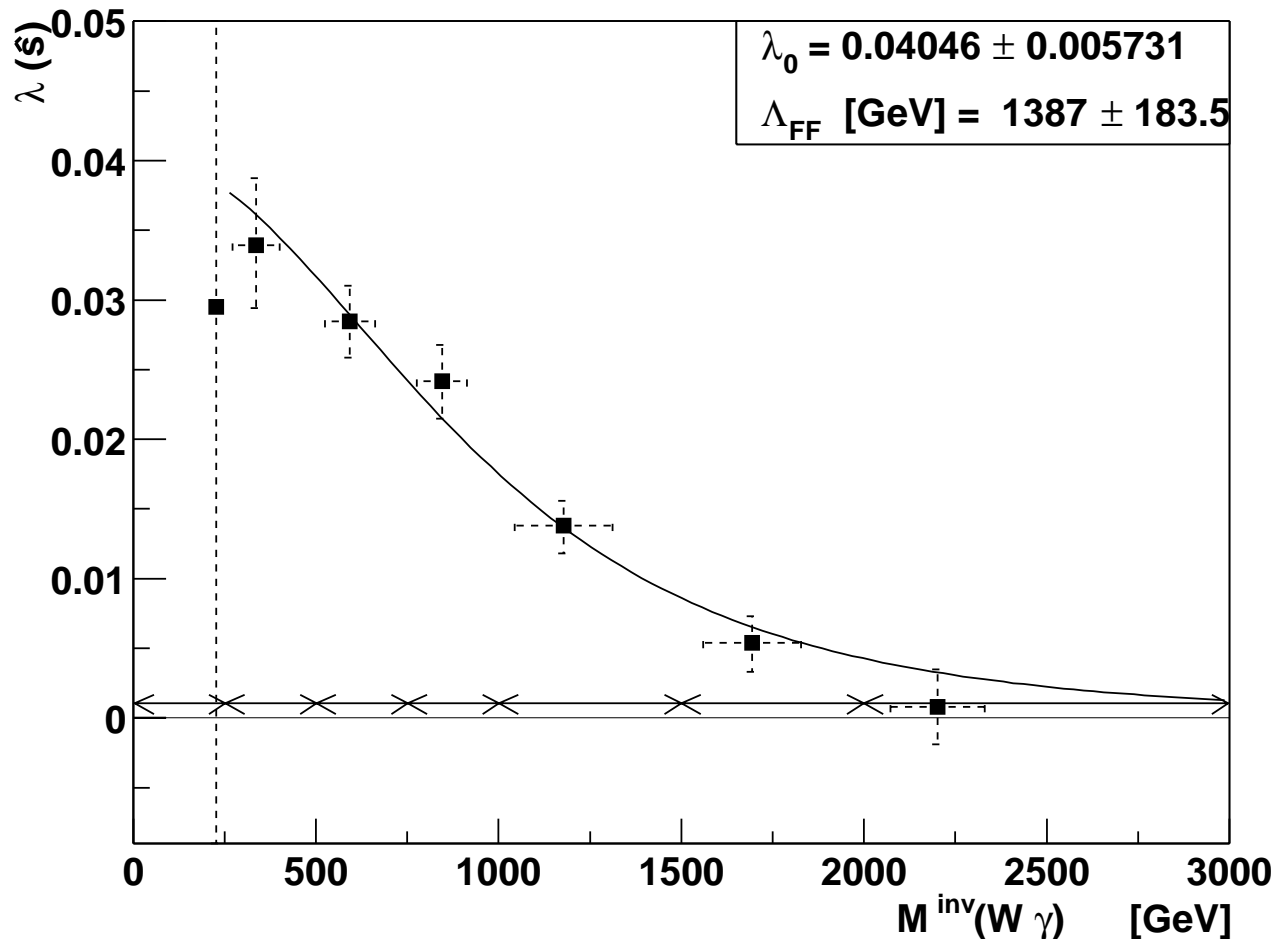


- Statistics will dominate LHC measurements (except for Δg^1)
→ sensitivity derived from a few events in the high $P_T(V)$ tail
- Dominant systematics are theoretical:
→ neglected higher orders and pdf's
- Systematics reported here are worst case scenario,
→ assumes we are unable to correct for the mis-modeling.

What can we do if we observe anomalous couplings??



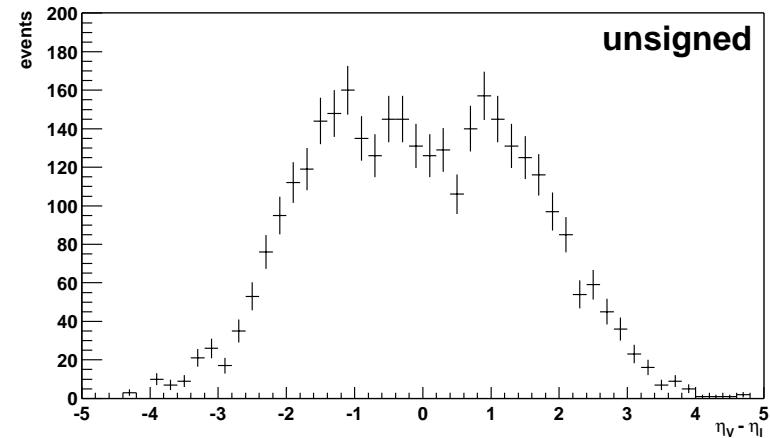
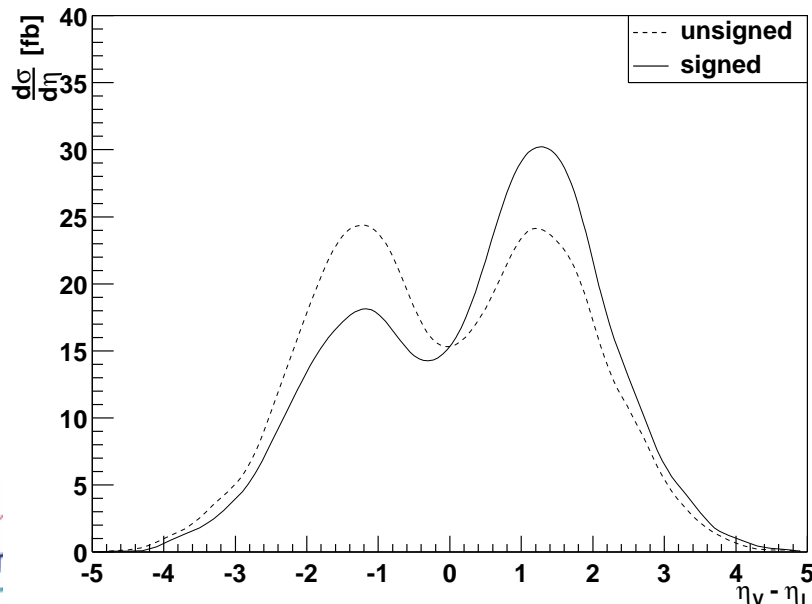
→ LHC will have sufficient statistics to measure the form factor behavior (energy dependence) of anomalous couplings.





- specific production angle of the photon is forbidden by subtle gauge cancellations
 - one of very few remaining electroweak discoveries
 - normal statement "Tevatron has a distinct advantage because of the asymmetric beams."
- ➔ borrow the (Drell-Yan) idea of signing forward direction by the system boost.

Radiation Zero





Triple Gauge-boson Couplings:

RESULTS & Summary

- 95% Confidence Intervals are:
 - limits derived by averaging over 5000 "mock" ATLAS expts.
 - typically **order of magnitude improvement** over LEP / Tevatron
 - statistically limited measurement (!)
 - ⊗ **sensitivity from a few events in high P_T tail** (except Δg_Z^1 , for which systematics & statistics are comparable)
 - theoretical errors dominate the systematics
 - ⊗ "tools" for controlling these systematics have been developed, not discussed here.
 - new means of ensuring unitarity developed (not discussed here)
 - measurements of anomalous couplings as a function of energy will be possible.
- $-0.0035 < \lambda_\gamma < +0.0035$**

$-0.0073 < \lambda_Z < +0.0073$

$-0.075 < \Delta\kappa_\gamma < +0.076$

$-0.11 < \Delta\kappa_Z < +0.12$

$-0.0086 < \Delta g_Z^1 < 0.011$
- For 30 fb⁻¹, systematics included.**



Conclusions

- ATLAS is under construction.
 - ⊗ performance requirements are being met in beam tests.
 - ⊗ physics studies drive the performance goals.
- ATLAS physics potential includes *competitive precision* electroweak measurements: $\sin^2\theta_W$, $\text{mass}(W)$, TGCs,...
- new Monte Carlo techniques for combining NLO(α_s) matrix elements with the parton shower approach have been developed → excellent **tool for** (by!) **experimentalists**.
- Triple Gauge-boson couplings probe the very foundation of the Standard Model.
 - measurements will be statistically limited, even at LHC
 - order of magnitude improvement in confidence limits over previous expts.
 - new means of ensuring unitarity (form factors) has been introduced (not discussed in this talk, ask if you're interested)

